

Multiple stellar populations in star clusters

Or:

Gluttonous Star Clusters: when they can not get enough

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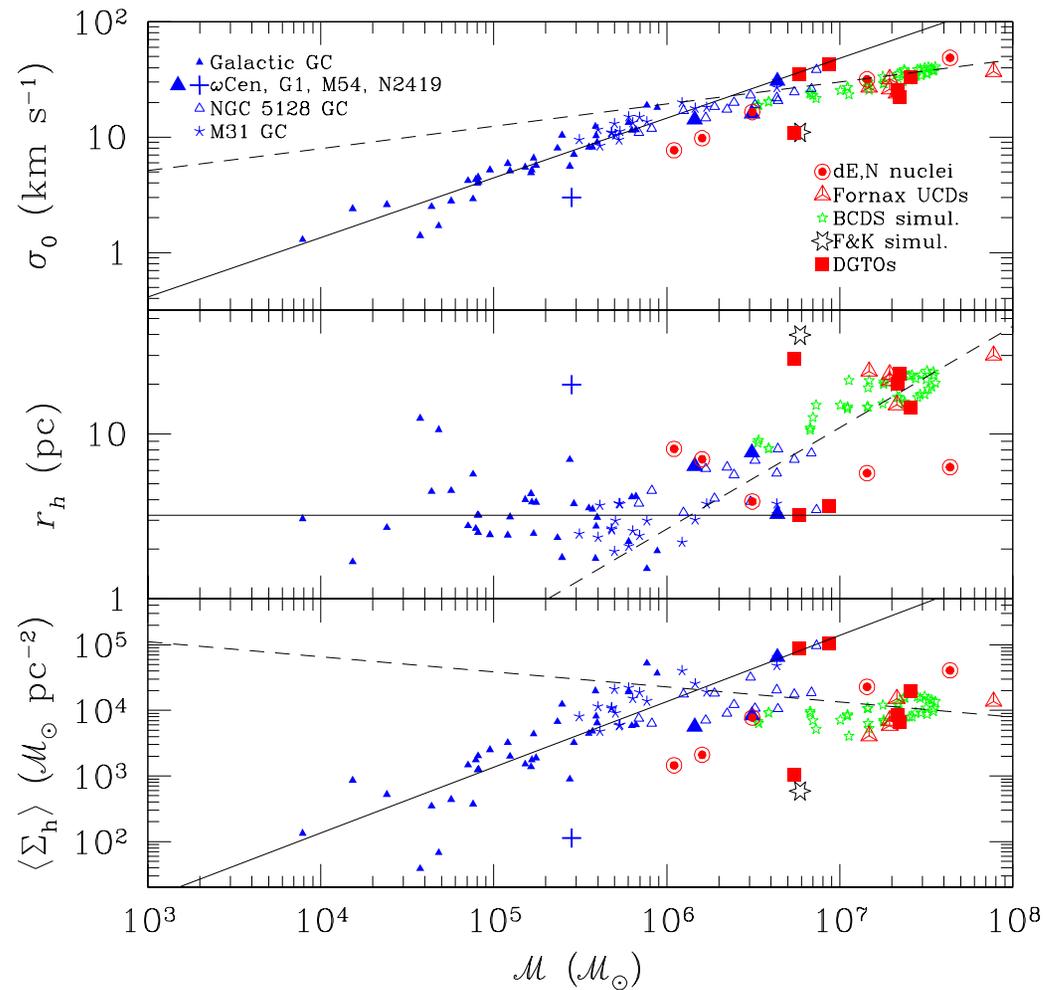
Sternwarte Bonn

Peculiar clusters

cluster	M_{cl}	peculiarity	reference
Ω Cen	$2.5 \cdot 10^6 M_{\odot}$	metallicity spread three stellar populations	Freeman & Rodgers (1975) Hilker & Richtler (2000)
G1	$8 \cdot 10^6 M_{\odot}$	metallicity spread	Meylan et al. (2001)
M54	$2 \cdot 10^6 M_{\odot}$	metallicity spread	Sarajedini & Layden (1995)
NGC 6388	$1.1 \cdot 10^6 M_{\odot}$	two distinct populations?	Ree et al. (2002)
NGC 6441	$1.6 \cdot 10^6 M_{\odot}$	two distinct populations?	Ree et al. (2002)

masses are from McLaughlin & van der Marel (2005), G1 from Baumgardt et al. (2003)

Mass-Radius-Relation



Haşegan et al. (2005)

Nuclear clusters

Walcher et al. (2005):

- 9 cluster in late-type bulgeless spiral galaxies
- $M_{cl} = 8 \cdot 10^5 - 6 \cdot 10^7 M_{\odot}$
- $M/L_I^{dyn} = 0.17 - 1.46$

Suggestion: *"...NCs are protobulges that grow by repeated accretion of gas and subsequent star formation."*

What happens at

$10^6 M_{\odot}$?

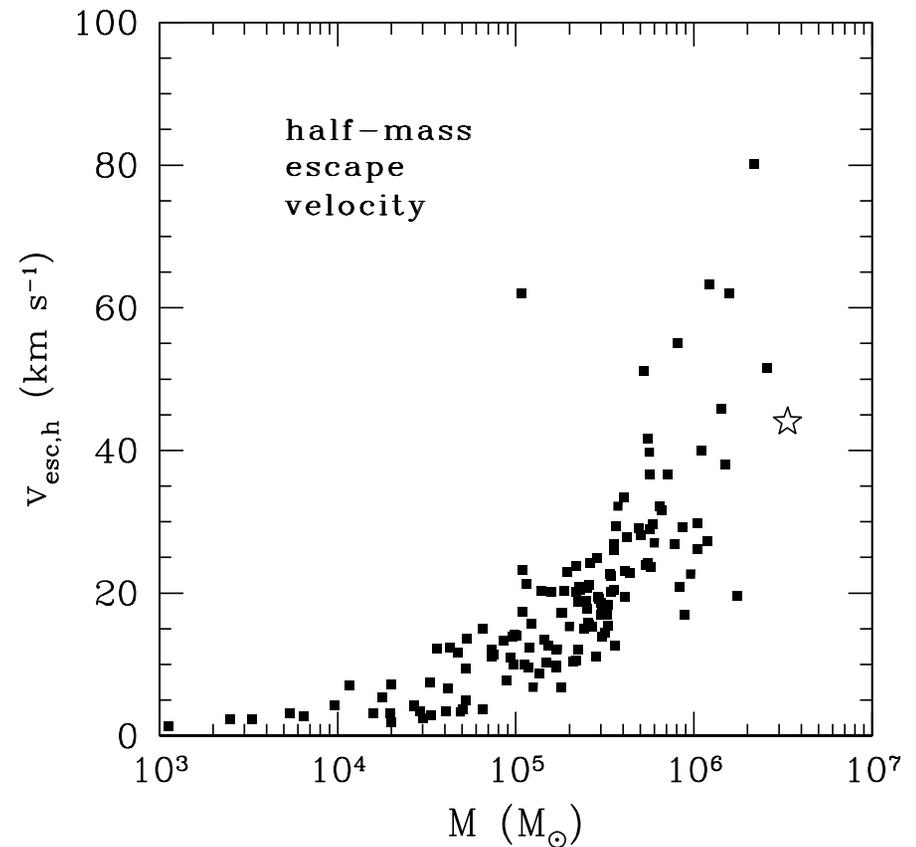
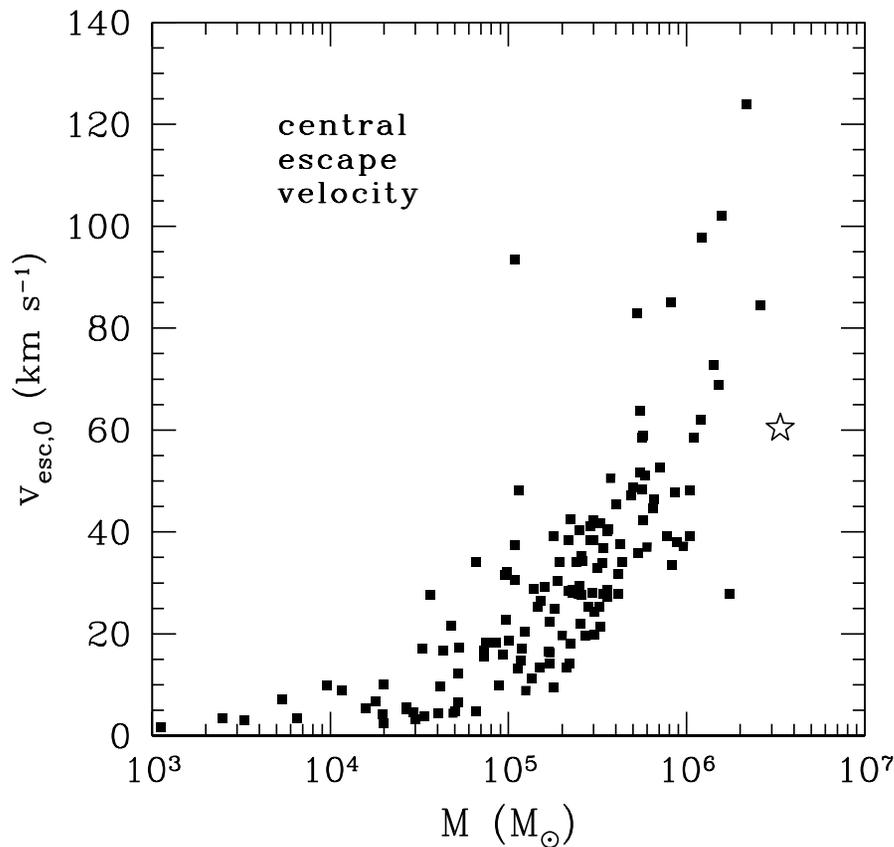
All peculiarities
are associated with
the observed or assumed
existence of
multiple stellar populations.

Material from inside?

- Very massive star clusters can retain AGB-wind material (Smith et al., 2000; D'Antona & Caloi, 2004).
- terminal velocity of AGB winds $\sim 15 \text{ km s}^{-1}$ (Loup et al., 1993)

Material from inside?

Escape velocity of 141 MW-GC



Gnedin et al. (2002)

Material from outside?

ISM components

phase	T/K	n/cm^{-3}	filling factor
molecular	10–20	$10^2\text{--}10^6$	1–2
cold atomic	50–100	20–50	
warm atomic	6000–10 000	0.2–0.5	~ 50
warm ionised	~ 8000	0.2–0.5	
hot ionised	$\sim 10^6$	~ 0.0065	~ 50

Ferrière (2001)

Model: basic equations

$$(1) \quad \frac{\partial \rho}{\partial t} + \nabla \bullet (\rho \mathbf{v}) = 0$$

$$(2) \quad \frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \bullet \nabla) \mathbf{v} = -\frac{\nabla P}{\rho} - \nabla \Phi$$

$$(3) \quad \frac{\partial \epsilon}{\partial t} + \mathbf{v} \bullet \nabla \epsilon = -P \nabla \bullet \mathbf{v} + \frac{\Gamma}{\rho} - \frac{\Lambda}{\rho}$$

$$(4) \quad \Delta \Phi = 4\pi G \rho$$

$$(5) \quad P = f(\rho, \epsilon)$$

Model: simplified

static & stationary case:

$$(6) \quad \frac{1}{\rho} \nabla P = -\nabla \Phi$$

only external potential:

$$(7) \quad \Phi(r) = -GM_{\text{cl}} (r^2 + b_{\text{pl}}^2)^{-1/2}$$

isothermal ideal gas

$$(8) \quad P = \rho \frac{k_{\text{B}}}{\mu m_{\text{u}}} T$$

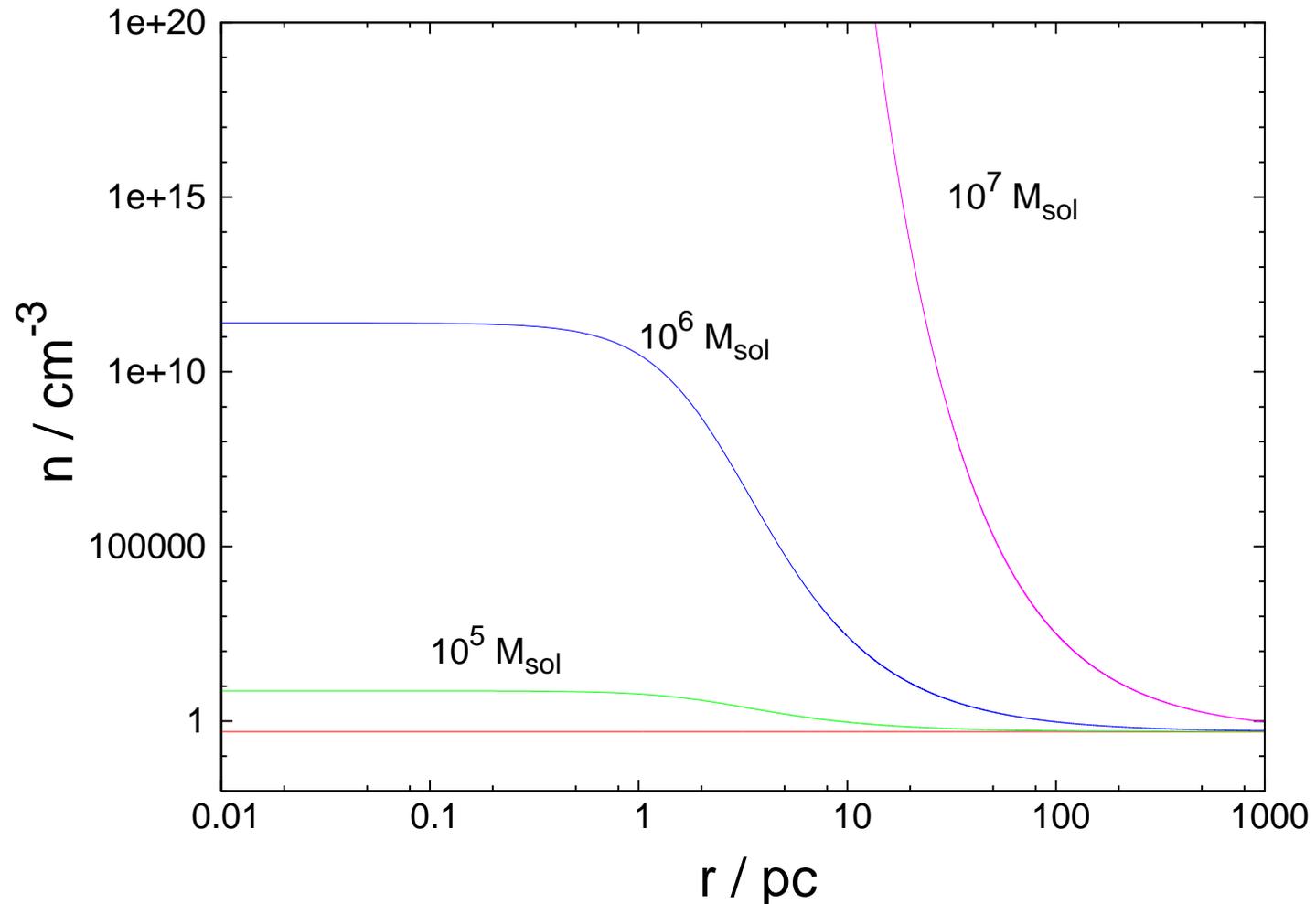
Model: solution

isothermal gas in equilibrium:

$$(9) \quad \rho = \rho_0 e^{-\frac{\mu m_{\text{u}}}{k_{\text{B}} T} (\Phi(r) - \Phi_0)}$$

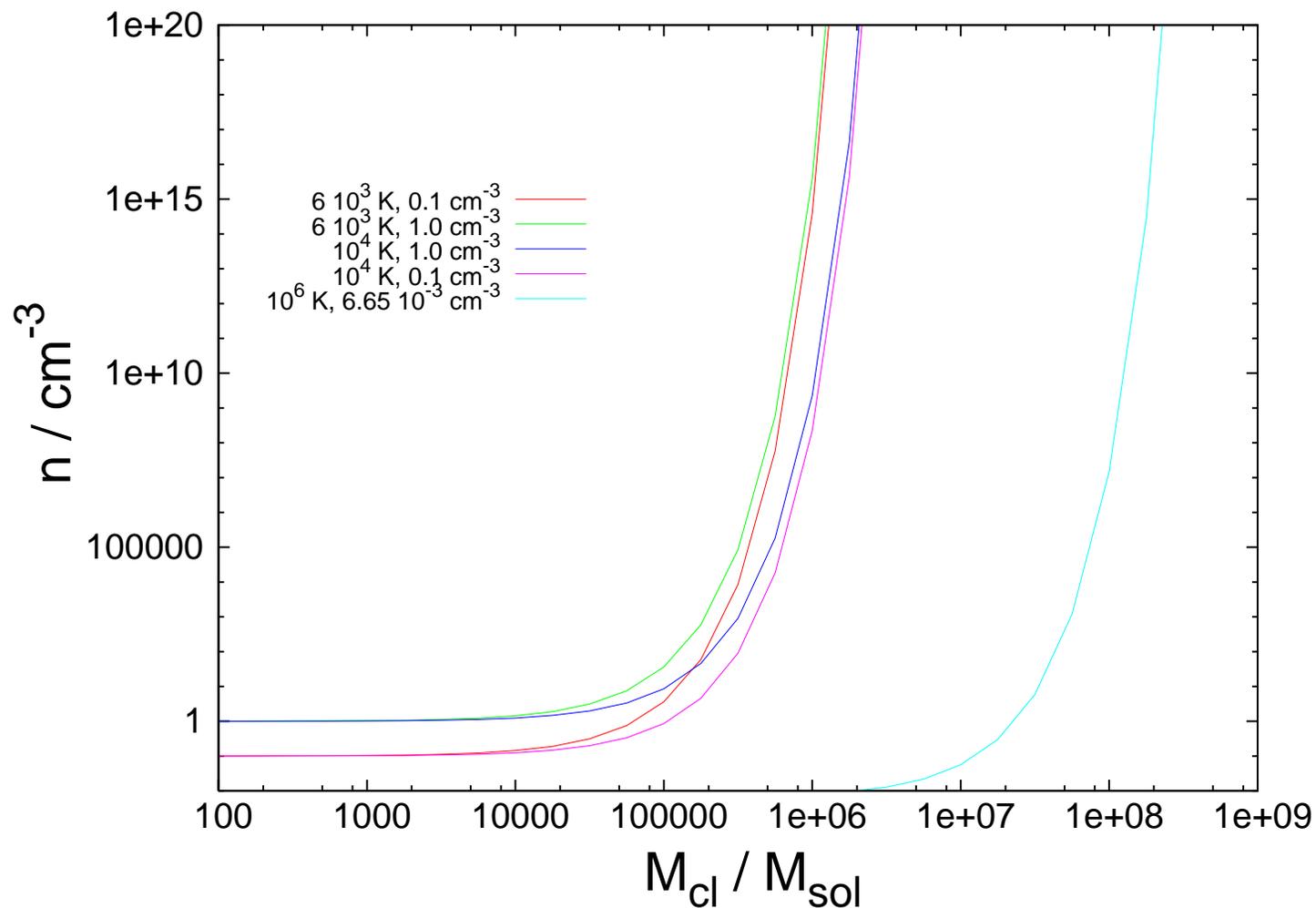
at $+\infty$: $\Phi_0 = 0$, $\rho_0 = \rho_{\text{ISM}}$

Density profile



warm ISM: $T=8000 \text{ K}$, $n_0=0.5 \text{ cm}^{-3}$
cluster: $\bar{r}_h=3.2 \text{ pc} \leftrightarrow \bar{b}_{\text{pl}}=2.4 \text{ pc}$

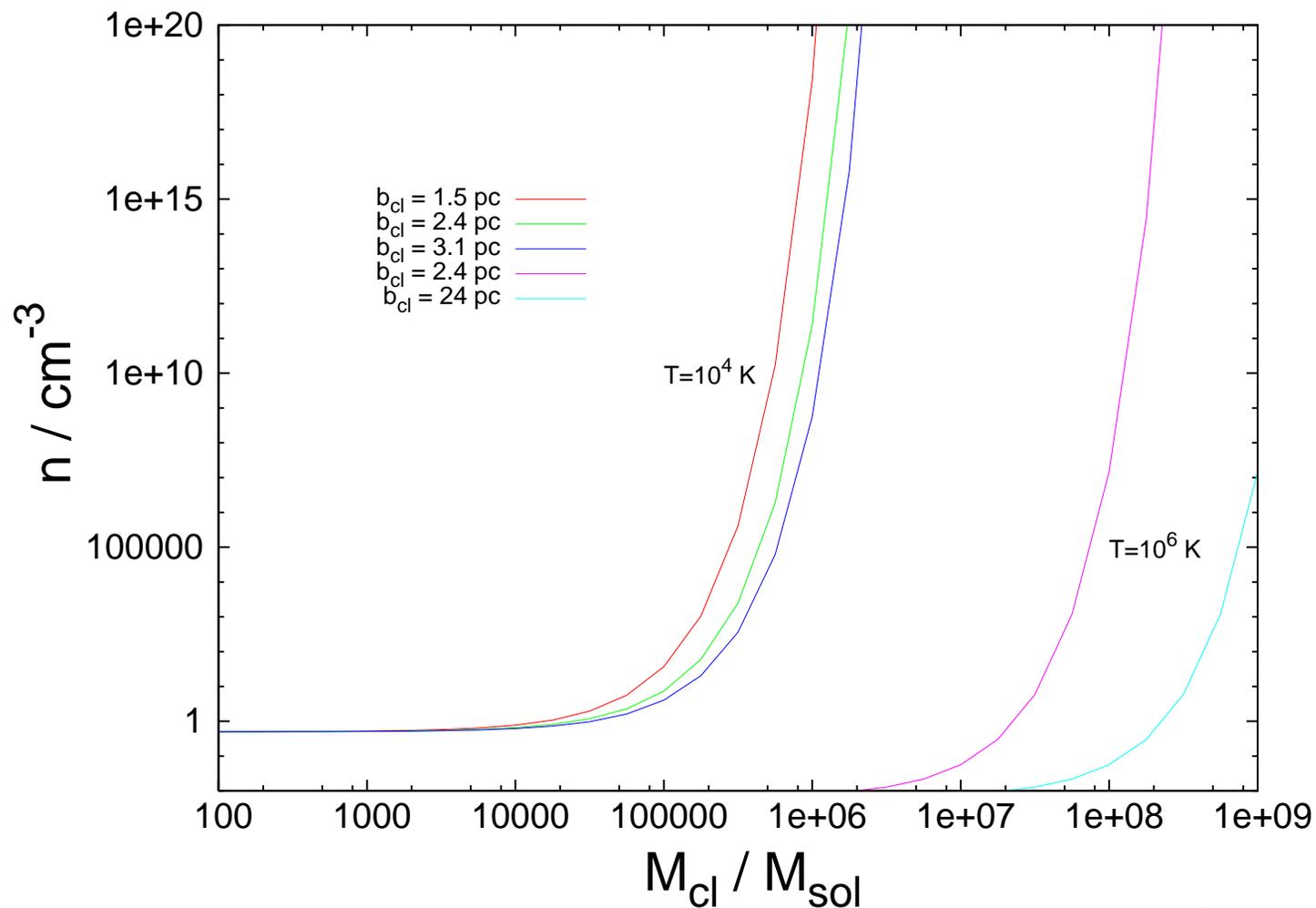
Central particle density – cluster mass (1)



warm ISM: $T=6000-10\,000$ K, $n_0=0.1-1.0$ cm^{-3}

cluster: $\bar{r}_h = 3.2$ pc \leftrightarrow $\bar{b}_{\text{pl}}=2.4$ pc

Central particle density – cluster mass (2)



warm ISM: $T = 8000 \text{ K}$, $n_0 = 0.5 \text{ cm}^{-3}$
cluster: $\bar{r}_h = 2.0 - 4.2 \text{ pc} \leftrightarrow \bar{b}_{\text{pl}} = 1.5 - 3.1 \text{ pc}$

Conclusions

- In the static isothermal model the density of the warm component of the ISM becomes singular at the cluster centre above a cluster mass of $10^6 M_{\odot}$ and for observed cluster half-light radii.
- Runaway cooling leads to further accretion of gas until stopped by star formation.
- Massive stars cluster are condensation germs for (giant) HI/molecular clouds: both within them or behind them.
- Increasing mass of a cluster leads to a faster orbit decay through dynamical friction.

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