DYNAMICS OF ROTATING CLUSTERS WITH BLACK HOLES

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Black Holes in Globular Clusters - Motivation

Black Holes in globular clusters:

Gebhardt et al. (2000), M15 Gerssen et al. (2002, 2003), M15 Gebhardt et al. (2002), G1 (M31)



NASA (AUR/STScI) and M. Rich (UCLA)



Chandra image of a dense cluster in M82

Rotation in globular clusters:

Meylan & Mayor (1986), ω Cen , 47 Tuc Lupton, Gunn & Griffin (1987), M13 Van Leeuwen (2000), ω Cen Anderson & King (2003), 47 Tuc McLaughlin et al. (2005), 47 Tuc van de Ven et al. (2006), ω Cen



 $e=1 - b/a \sim 0.1$ (MW GCs) to 0.2 (LMC GCs)

Black Holes in Globular Clusters - Motivation

Numerical models :

<u>Direct N-Body</u>: individual orbits of high N are integrated Spurzem & Aarseth (1996), Baumgardt & Makino (2003), Ardi et al. (2005)*, Baumgardt et al. (2004)

<u>Monte Carlo</u>: diffusion of (random generated) velocities are calculated at selected test particles

Marchant & Shapiro (1980), Spurzem & Giersz (1996), Giersz & Spurzem (2003)

Fokker Planck: Boltzmann equation + collisional term

Lightman & Shapiro (1977), Takahashi (1995,1996,1997), Einsel & Spurzem (1999)*, Kim et al. (2002, 2004)*

<u>Gaseous Model:</u> Momentum of the Fokker Planck equation Louis & Spurzem (1991), Amaro-Seoane et al. (2004) (* rotation)

,Seed' BHs in GCs:

Portegies Zwart et al. (2002) Baumgardt et al. (2005) Freitag et al. (2005)

<u>Aims:</u>

- Implications in evolution of rotation in relaxed systems with BH

- Link between stellar-mass and supermassive BHs

<u>Fokker Planck model</u> <u>Initial conditions</u>:

- single-mass system axisymmetric in space, anisotropic in velocity space
- cluster evolution time scales:

$$au_{
m cl} > au_{
m rh} \gg au_{
m dyn}$$

$$au_{
m rh} = rac{0.138 \sqrt{N r_{
m h}^3}}{\sqrt{Gm} \ ln\Delta}$$

• Changes in f(x,v) due to small angle scatterings

• *M*_{*} < *M*_{bh} << *M*_{cl}

• Tidal galaxy boundary

$$m = (4)^3 = M = (4)$$

 \vec{v} to $\vec{v} + \Delta \vec{v}$

$$\frac{r_{\rm tid}(t)^3}{r_{\rm G}^3} = \frac{M_{\rm cl}(t)}{M_{\rm G}}$$

• stars are **possibly** disrupted if
$$J_z <$$

leading to BH growing

 $\Delta v/v \ll 1$



Rotating King models:

$$f(E, J_z) \propto \exp\left(-\beta \Omega_0 J_z\right) \cdot \left[\exp\left(-\beta E\right) - 1\right]$$

	Model name	W_0	ω_0
$\mathbf{W} = \partial (A + A) - (2 - C - 0)$			
$W_0 = -\beta(\phi_c - \phi_t)$ (3, 0, 9)	A1	3.0	0.00
	A2	3.0	0.30
$\sqrt{2/(4-C-)}$ (2.2.2.2.2.2.2.2.1.2)	A3	3.0	0.60
$\omega_0 = \sqrt{9/(4\pi G n_c)} \Omega_0 (0.0, 0.3, 0.6, 0.9, 1.2)$	A4	3.0	0.90
	A5	3.0	1.20
	B1	6.0	0.00
	B2	6.0	0.30
	B3	6.0	0.60
	B4	6.0	0.90
$E = \frac{1}{2}v^2 + \phi_{al}(\overline{\omega}, z) + \phi_{bb}(\overline{\omega}, z)$	B5	6.0	1.20
$\frac{2}{2} = 2^{\circ} + \varphi_{\mathcal{U}}(\mathbf{\omega}, \mathbf{\omega}) + \varphi_{\mathcal{U}}(\mathbf{\omega}, \mathbf{\omega}) = 2^{\circ}$	C1	9.0	0.20
	C2	9.0	0.30
$J_z=arpi ec v \hat e_{lpha}$	C3	9.0	0.40
	C4	9.0	0.50

Star accretion:

$$r_{\rm d} \sim r_* (M_{\rm bh}/m_*)^{1/3}$$

Are all stars in orbits of $J_z < J_{z,min}$ disrupted?

$$J_z^{min}(E) = r_d \sqrt{2(E - GM_{\rm bh}/r_d)}$$







Evolution of density profile:

-7/4 cusp (Bahcall & Wolf 1976, Lightman & Shapiro 1977) forms inside influence radius (squares)



2D evolution of density: Effects of rotation



Evolution of Lagrange Radii (dependence on initial rotation ω_0): Outer mass shells are faster depleted



Evolution of cluster Radii:

Core radius (r_c) falls up to collapse, Influence Radius (r_a) larger after collapse,



Evolution of cluster masses: M_{bh} stalls (dM_{bh}/dt has a maximum) at collapse. $M_{bh} \sim 0.01 M_{cl} M_{cl}$ drops faster in BH-model

Model name	$M_{\rm BH}^{\rm stall} \ (M_{\odot})$	$dM/dt_{\rm max}~(M_{\odot}/{\rm yr})$
A1	$2.1 \cdot 10^4$	$1.77 \cdot 10^{-3}$
A2	$2.7 \cdot 10^4$	$1.32 \cdot 10^{-3}$
A3	$3.0\cdot 10^3$	$1.21 \cdot 10^{-3}$
A4	$4.4 \cdot 10^3$	$2.83 \cdot 10^{-4}$
A5	$5.0\cdot 10^3$	$3.24 \cdot 10^{-3}$
B1	$4.8\cdot 10^4$	$3.40 \cdot 10^{-4}$
B2	$4.2 \cdot 10^{4}$	$4.74 \cdot 10^{-4}$
B3	$8.5 \cdot 10^4$	$1.39 \cdot 10^{-3}$
B4	$9.5 \cdot 10^4$	$3.78 \cdot 10^{-3}$
B5	$4.5 \cdot 10^4$	$1.30 \cdot 10^{-2}$
C1	$5.0\cdot 10^3$	$5.38 \cdot 10^{-4}$
C2	$2.3 \cdot 10^3$	$1.87 \cdot 10^{-5}$
C3	$1.5 \cdot 10^3$	$4.29 \cdot 10^{-6}$
C4	$1.9 \cdot 10^3$	$4.76 \cdot 10^{-6}$

Evolution of rotational velocity : Gravogyro instabilities carry out angular momentum and core rotates faster. BH supports rotation



$$\delta\Omega\approx-\frac{\delta j}{r^2}$$

Inagaki & Hachisu (1978)



Distribution function f(E,Jz)

 $(-1 \leq Y \leq +1)$

Rotation dominates collapse



Set of models:

Ellipticity / concentration vs. time



Fiestas et al. (2006)

Comparison with observations:

Rotation in the Meridional plane (Meylan & Mayor, 1986) (Merritt et al., 1997)











Specific angular momentum:

NO BH:

BH:





Black Holes in Globular Clusters - Conclusions

Gravogyro + Gravothermal effects drive collapse

Acceleration of evolution due to rotation + BH accretion (faster collapse), and faster mass loss (shortening of life time)

Post-collapse driven by BH energy source

Equilibrium states in density and vel.disp. profiles are as in nonrotating models

Rotation grows in the core limited by angular momentum loss (continuously transported outwards) and BH-accretion efficiency of most low a.m. (radial) orbits.

Data set of models and high resolution in space allows comparison with future observations

Future work: multimass-model, stellar evolution, galaxy tidal field, comparison to N-Body models. Implications in BH mass/sigma correlation