

The 11th BONN workshop on neutron stars

Testing Strong-Field Gravity with Pulsars and GWs

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MAX-PLANCK-GESELLSCHAFT

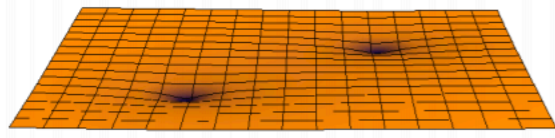
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Radioastronomie



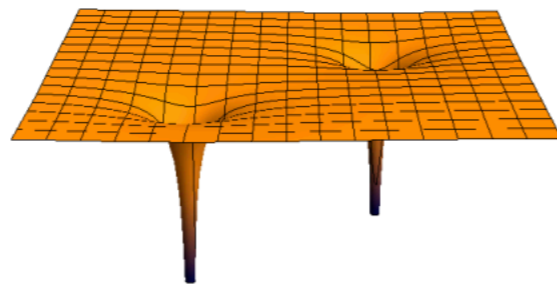
Outline

- ◆ Strong-field gravity (SFG)
- ◆ Testing SFG with pulsars
- ◆ Testing SFG with GW detectors
 - ✓ Binary black holes (BBHs): GW150914, GW151226, ...
 - ✓ Binary neutron stars (BNSs): GW170817
- ◆ Complementarity in testing SFG with pulsars and GWs

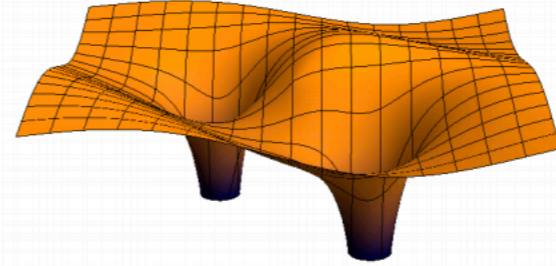
Strong-field gravity



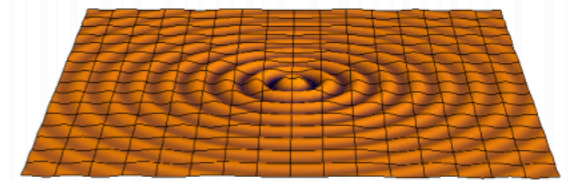
Solar System
G1



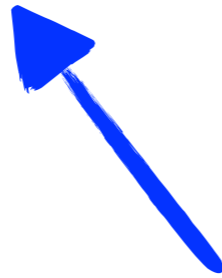
Binary Pulsar
G2



BBH Merger
G3



LIGO/Virgo Sites
GW



Strong-field gravity (SFG)

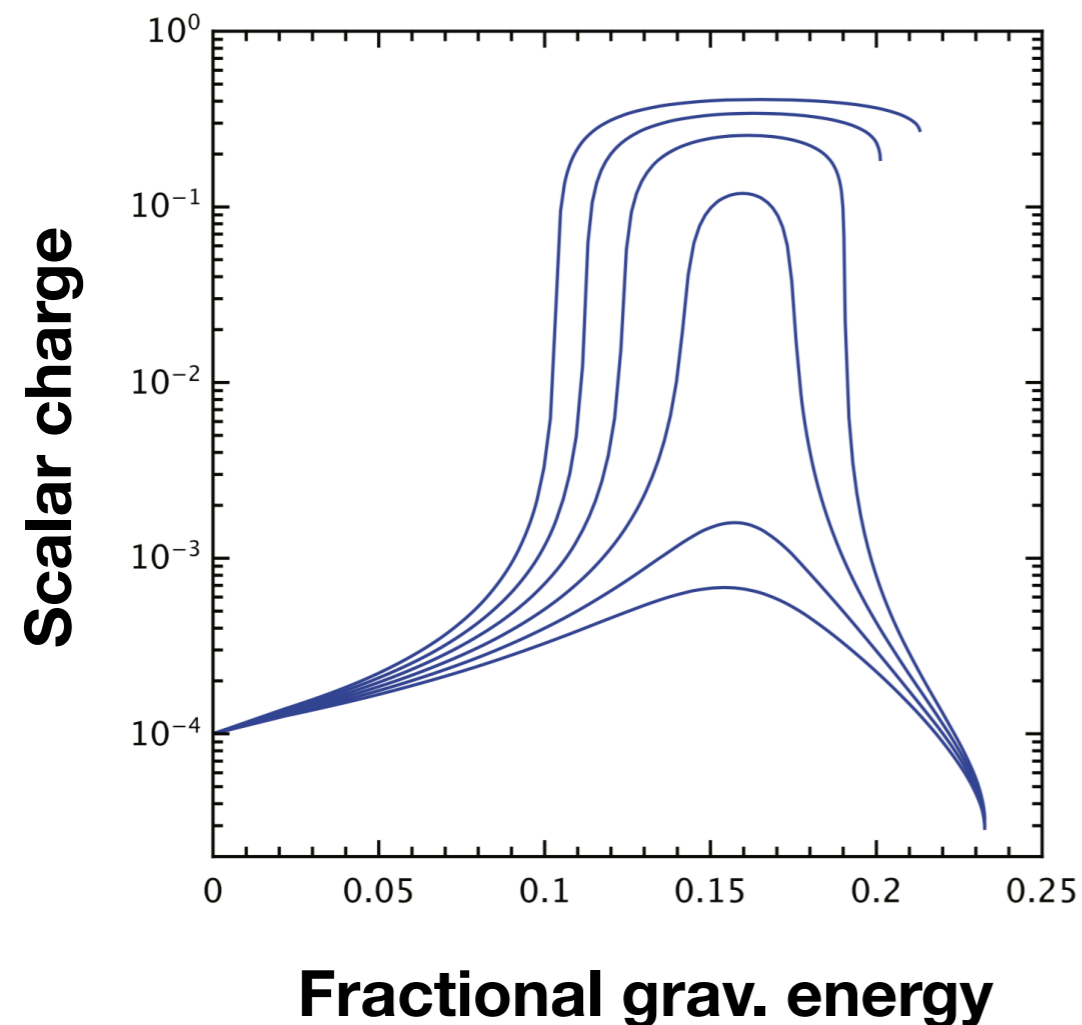
G1 – Quasi-stationary weak-field regime

G2 – Quasi-stationary strong-field regime

G3 – Highly dynamical strong-field regime

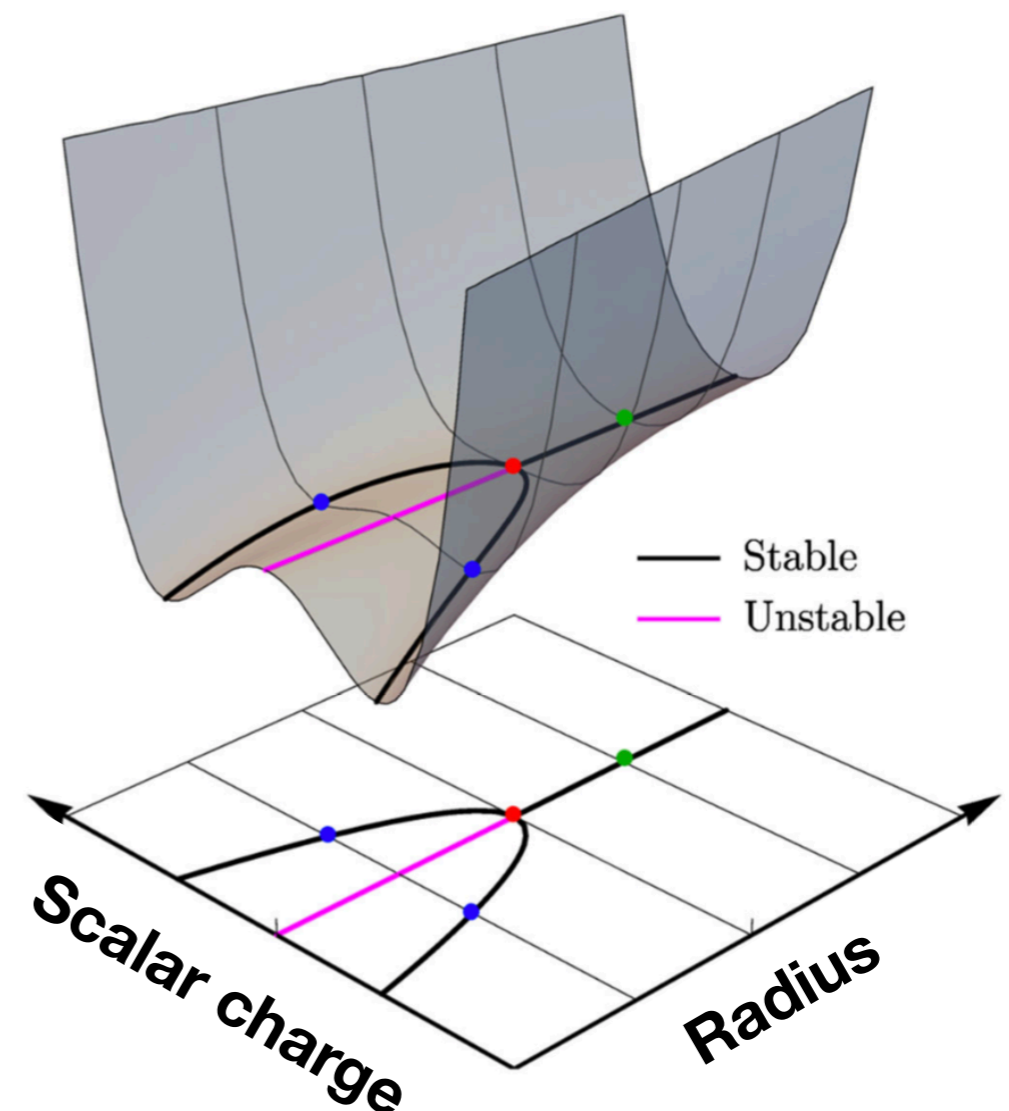
GW – Radiation regime

An example of SFG



Nonperturbative spontaneous scalarization could happen with NSs in a class of scalar-tensor theories

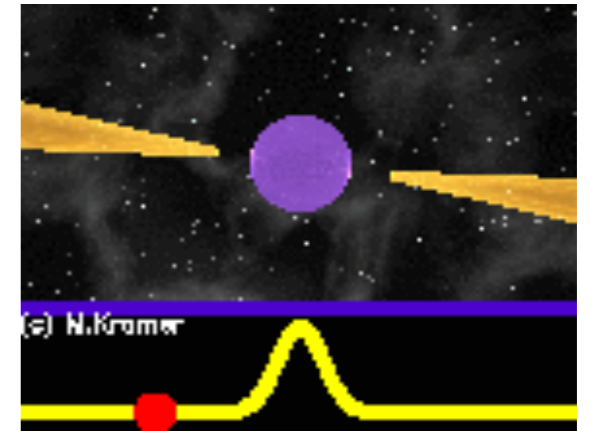
Damour & Esposito-Farèse 1993, PRL 70:2220
Damour & Esposito-Farèse 1996, PRD 54:1474
Shao & Wex 2016, SCPMA 59:699501



Strong-field behaviour is analogous to Landau's *phase transition* after a critical point

Damour & Esposito-Farèse 1996, PRD 54:1474
Esposito-Farèse 2004, AIP Conf. Proc. 736:35
Sennett, Shao, Steinhoff 2017, PRD 96:084019

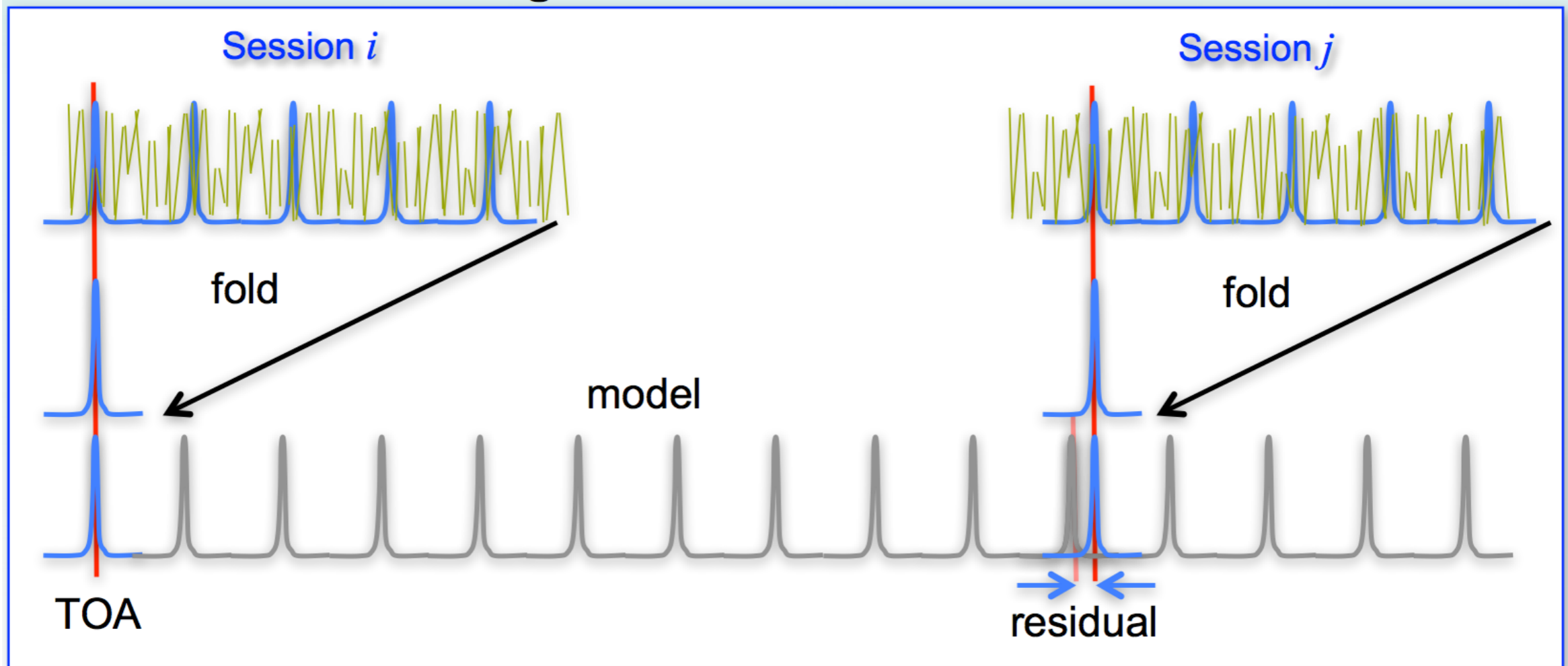
Pulsar timing



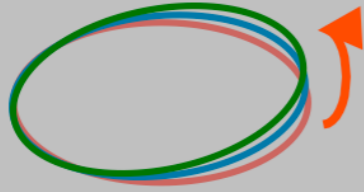
© M. Kramer

Phase-connected timing solution:

© N. Wex

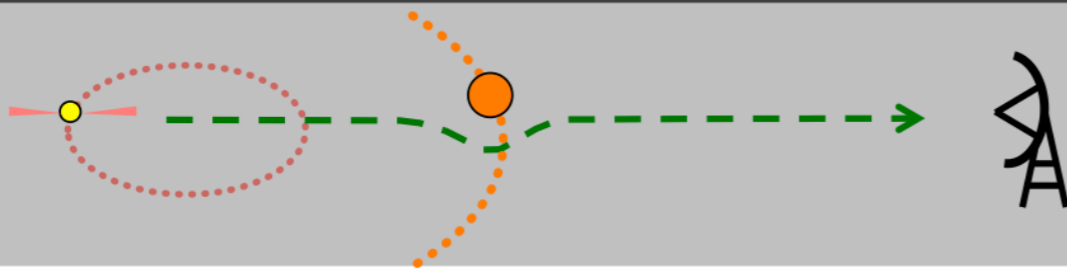


Timing of relativistic binary pulsars



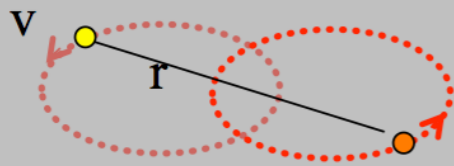
Precession

$$\dot{\omega} = 3 \frac{G^{2/3}}{c^2} \left(\frac{P_b}{2\pi} \right)^{-5/3} \frac{1}{1-e^2} \left[(m_1 + m_2) \right]^{2/3}$$



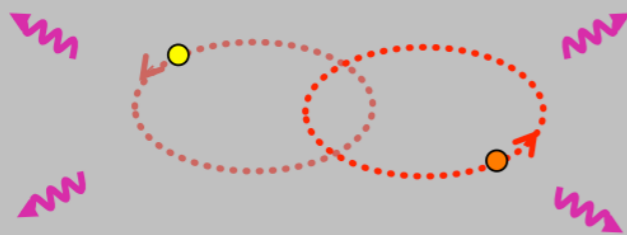
Shapiro Delay

$$\Delta t = 2 \frac{G}{c^3} m_2 \ln \left[\frac{1 + e \cos \varphi}{1 - \sin i \sin(\varphi - \varphi_0)} \right]$$



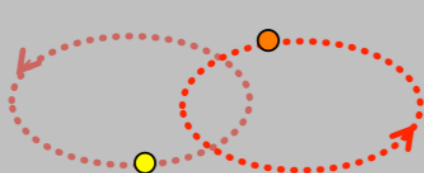
Grav Redshift/Time Dilation

$$\gamma = \frac{G^{2/3}}{c^2} \left(\frac{P_b}{2\pi} \right)^{1/3} e \frac{m_2 (m_1 + 2m_2)}{(m_1 + m_2)^{4/3}}$$



Gravitational Radiation

$$\dot{P}_b = - \left(\frac{192\pi}{5} \right) \frac{G^{5/3}}{c^5} \left(\frac{P_b}{2\pi} \right)^{-5/3} \left(1 + \frac{73}{24} e^2 + \frac{37}{96} e^4 \right) \frac{1}{(1-e^2)^{7/2}} \frac{m_1 m_2}{(m_1 + m_2)^{1/3}}$$



Second Orbit

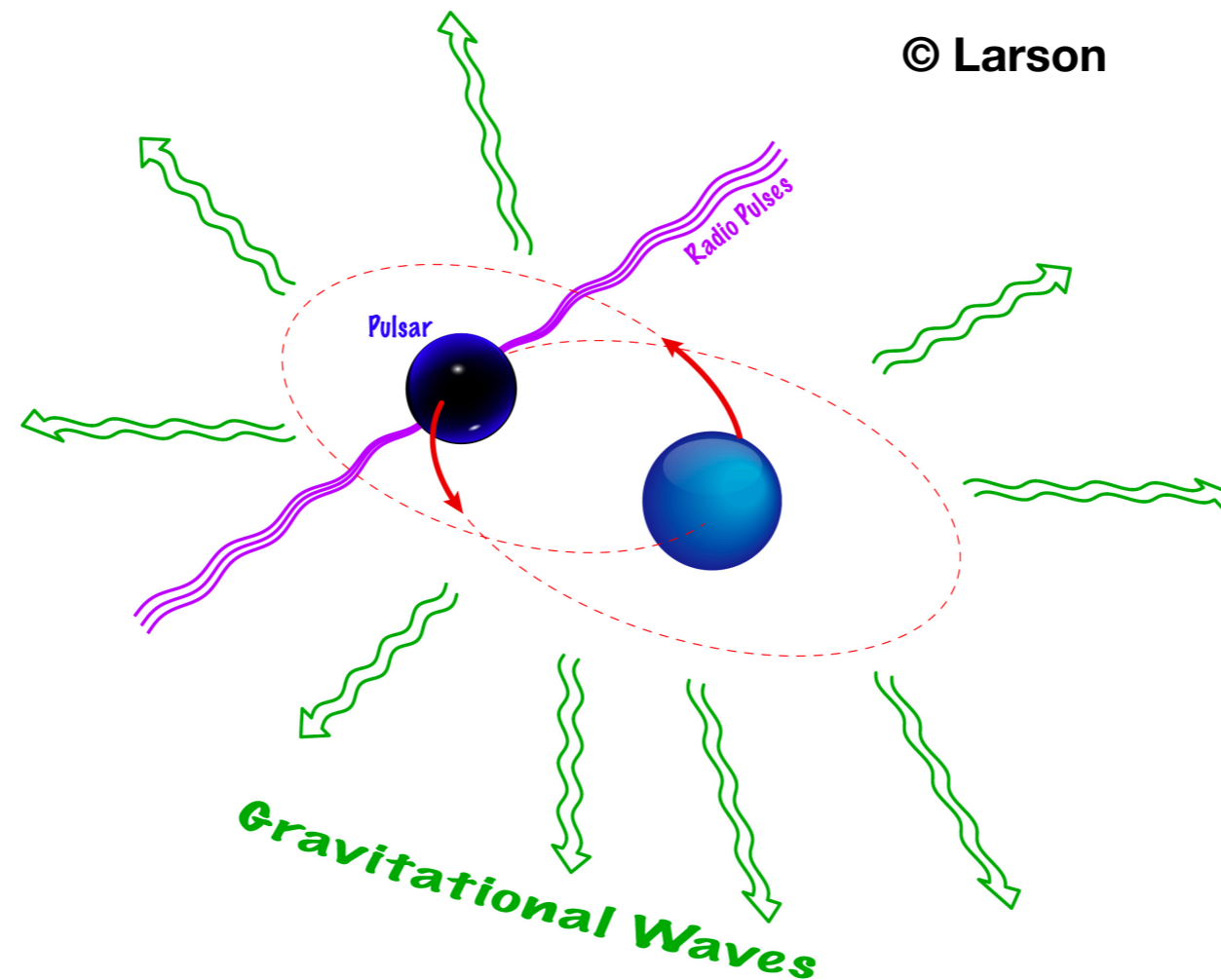
$$\frac{m_1}{m_2} = \frac{a_1 \sin i}{a_2 \sin i}$$



Physics 1993

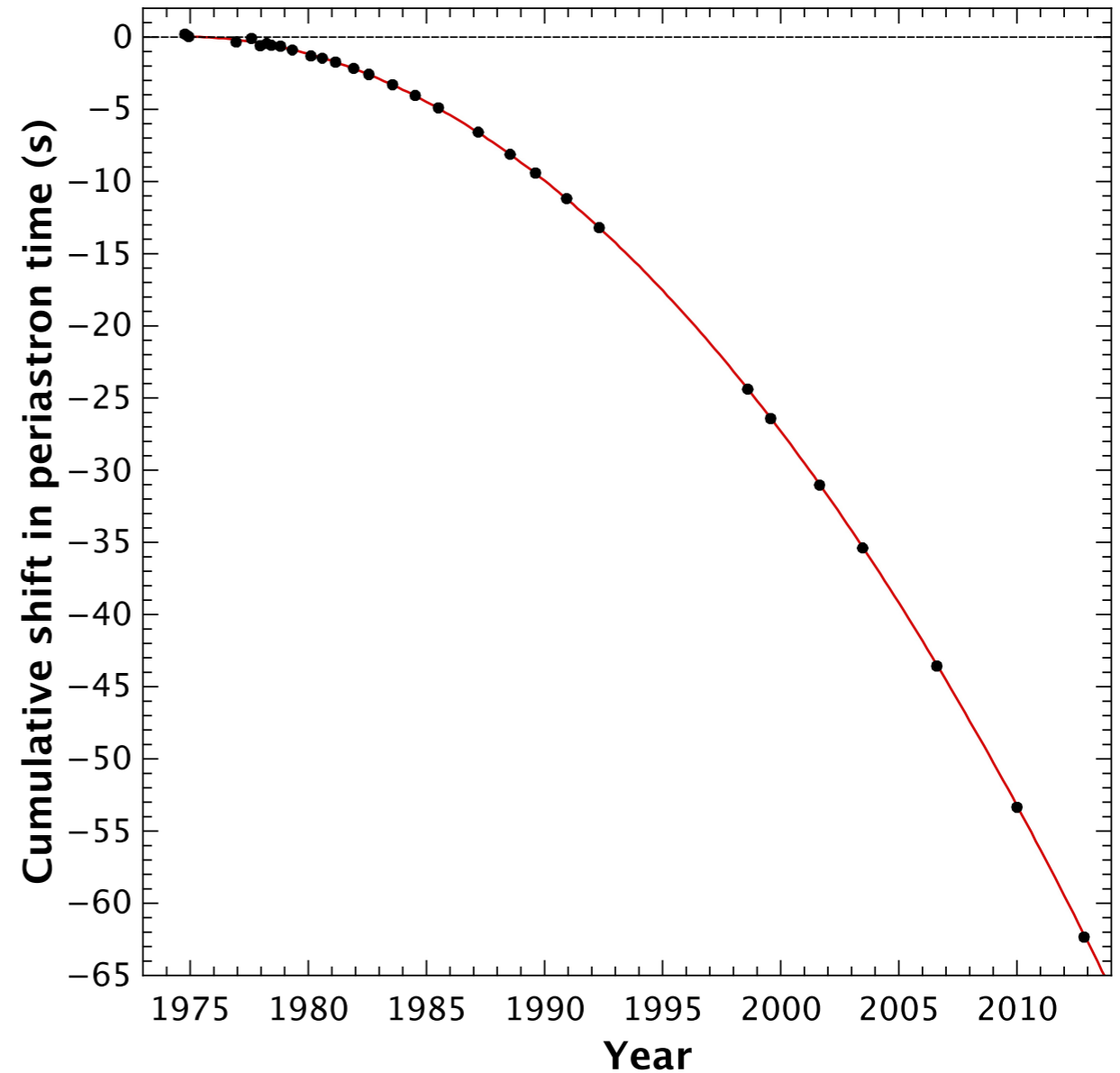
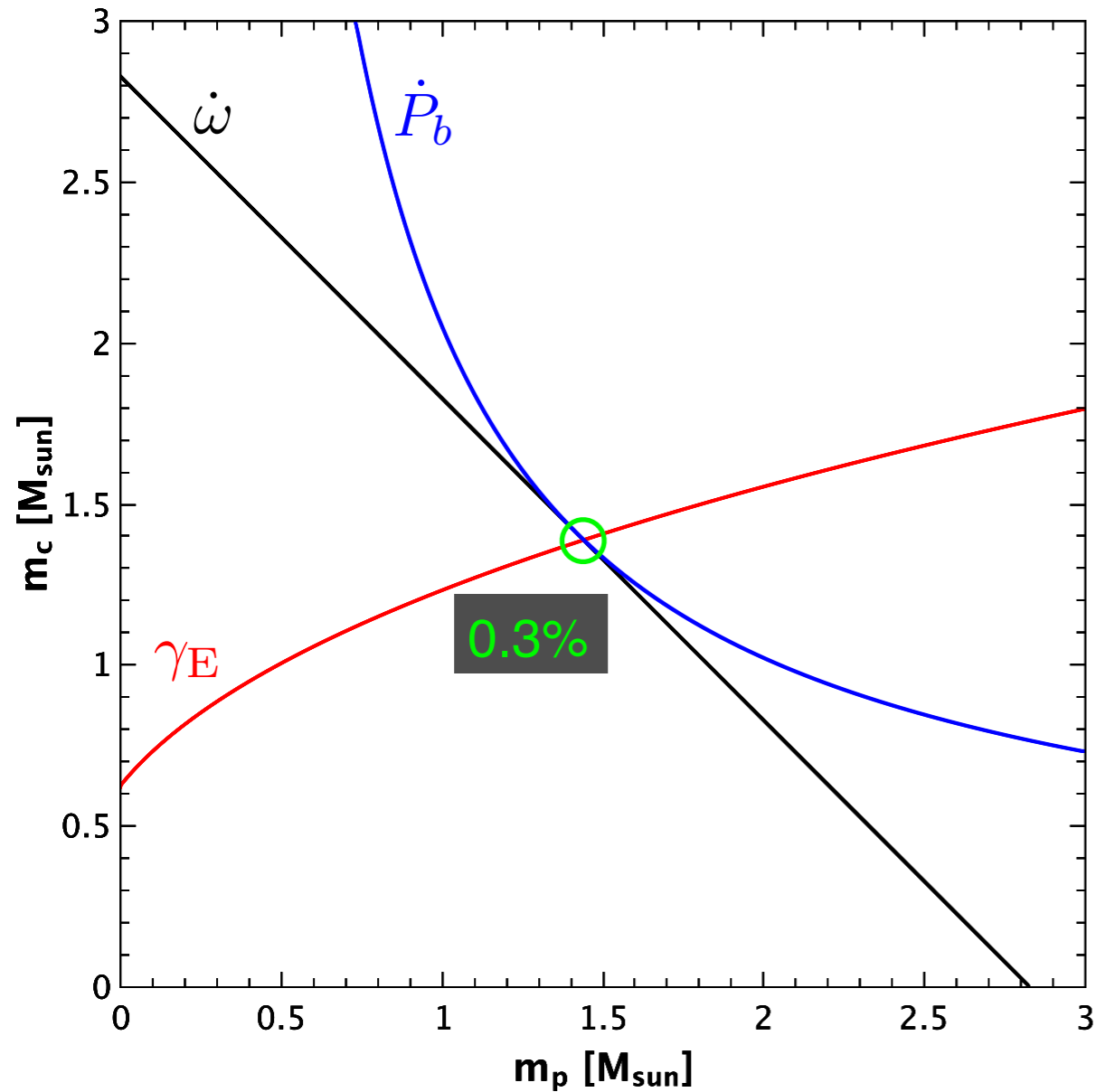
The Hulse-Taylor pulsar

Hulse & Taylor 1975, ApJ 195:L51



Radio pulsar timing provided the first evidence of GWs, and started a new era to test Einstein's gravity in strong field

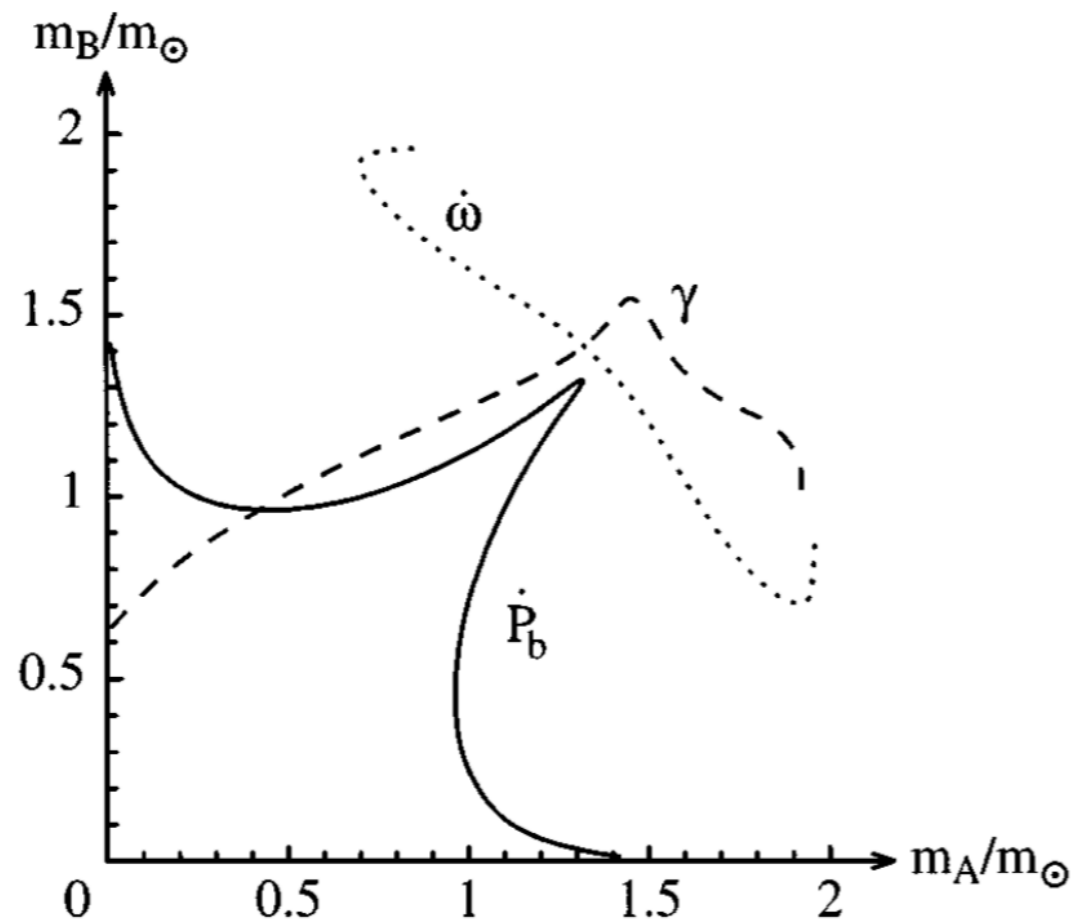
Mass-mass diagram and GW radiation



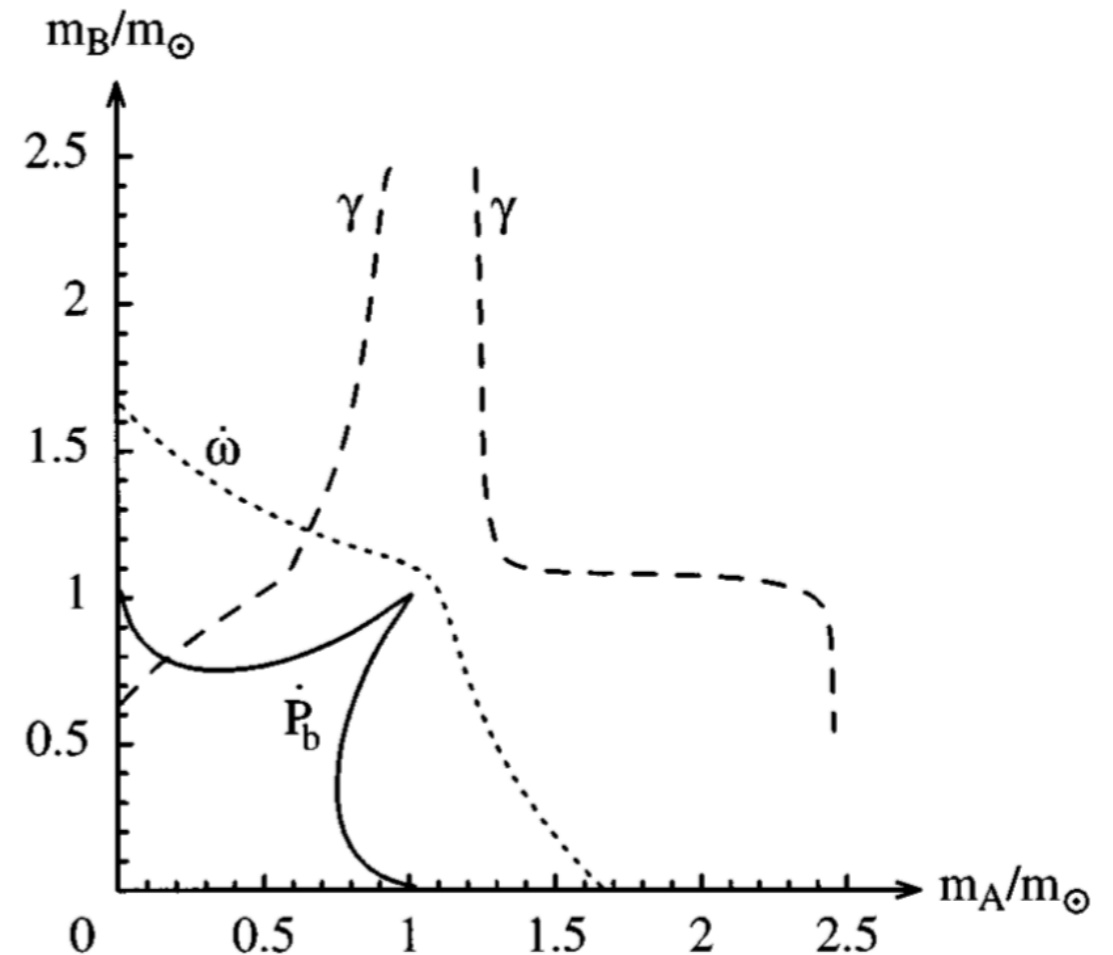
In GR, observations of post-Keplerian parameters agree within 0.3%

Weisberg & Huang 2016, ApJ 829:55

Examples: Hulse-Taylor pulsar with SFG



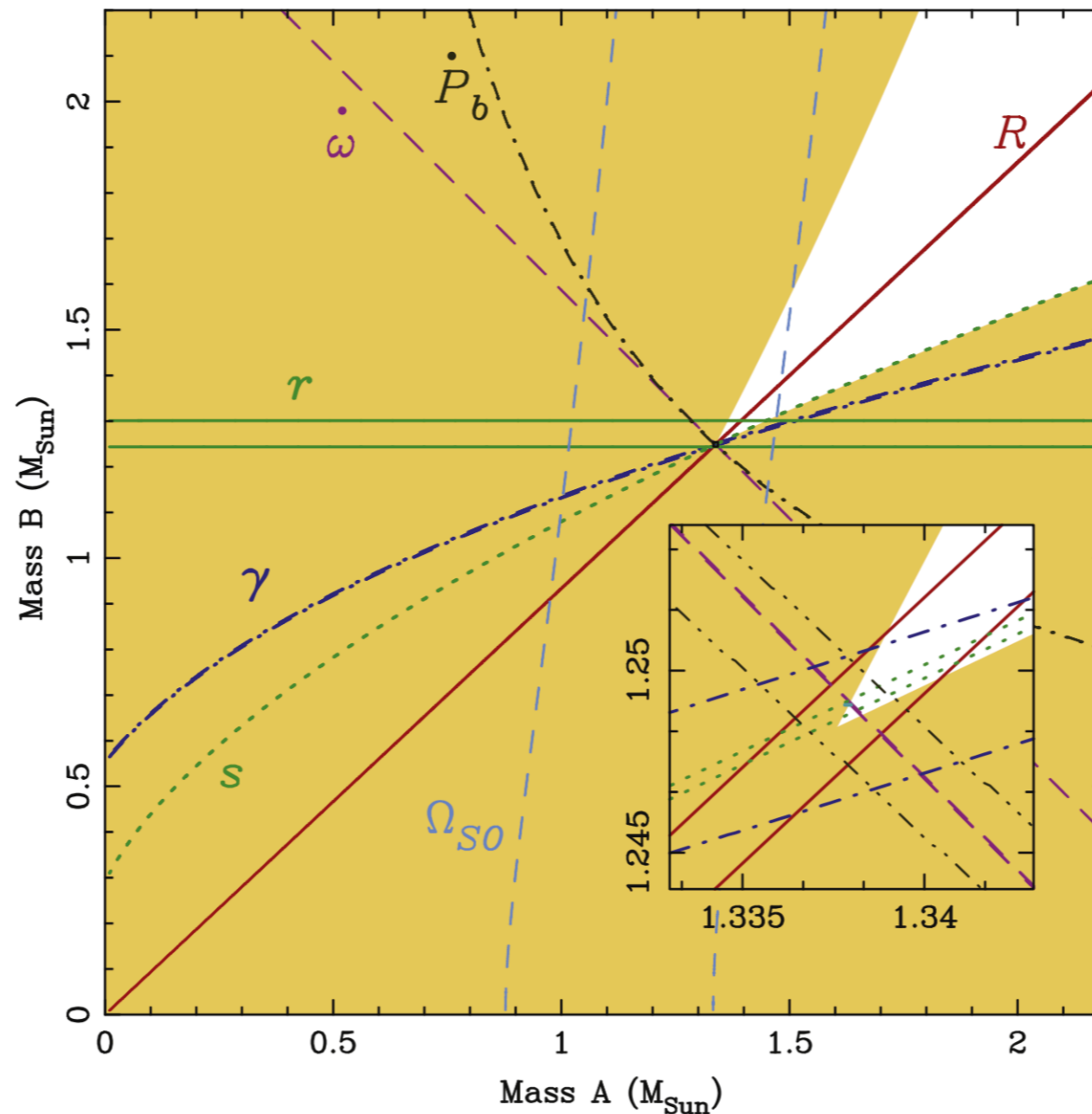
Scalar-tensor theory with $\beta_0 = -4.5$



Scalar-tensor theory with $\beta_0 = -6$

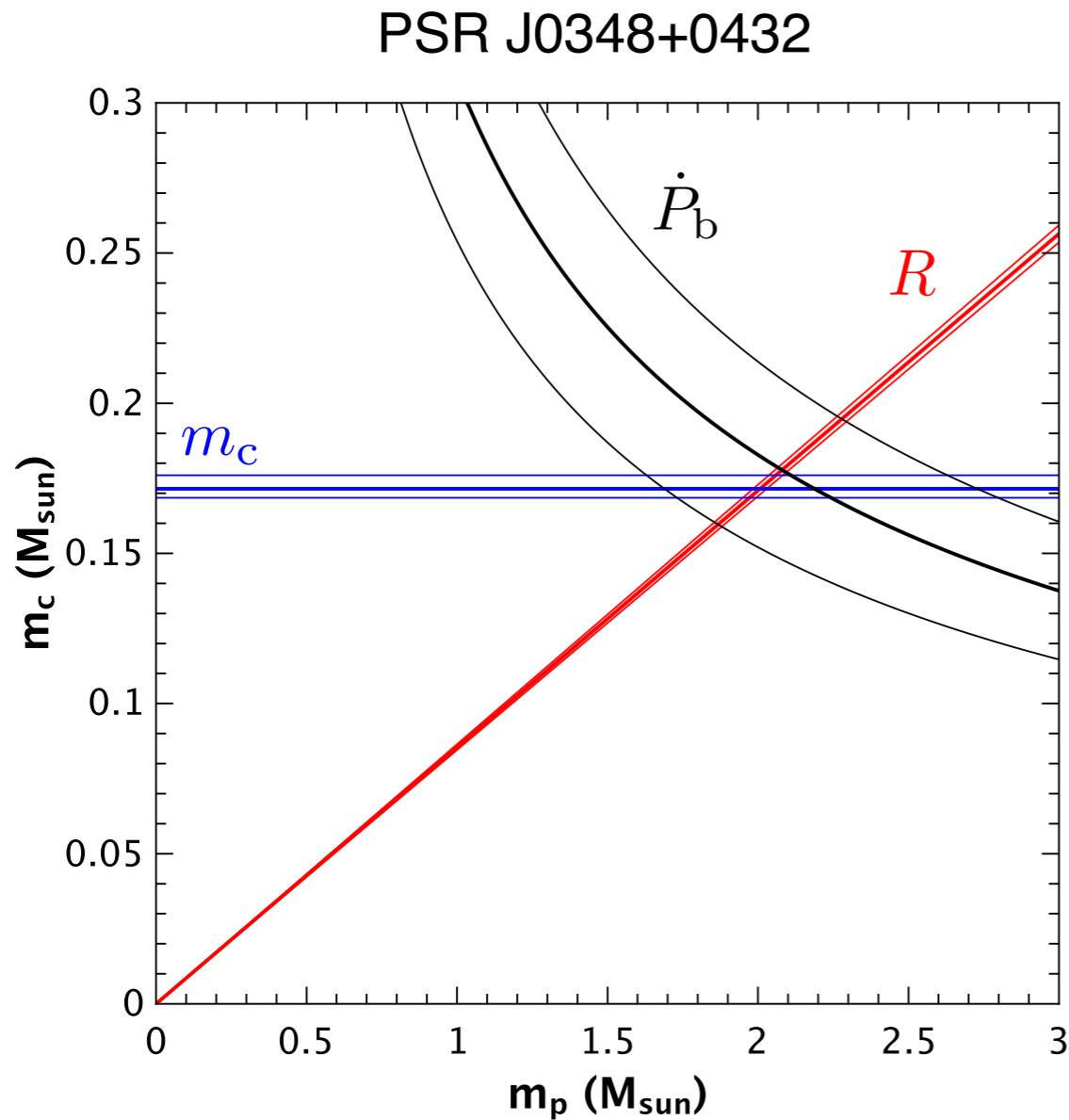
Binary pulsars are extremely sensitive probes of SFG!

The Double Pulsar J0737-3039

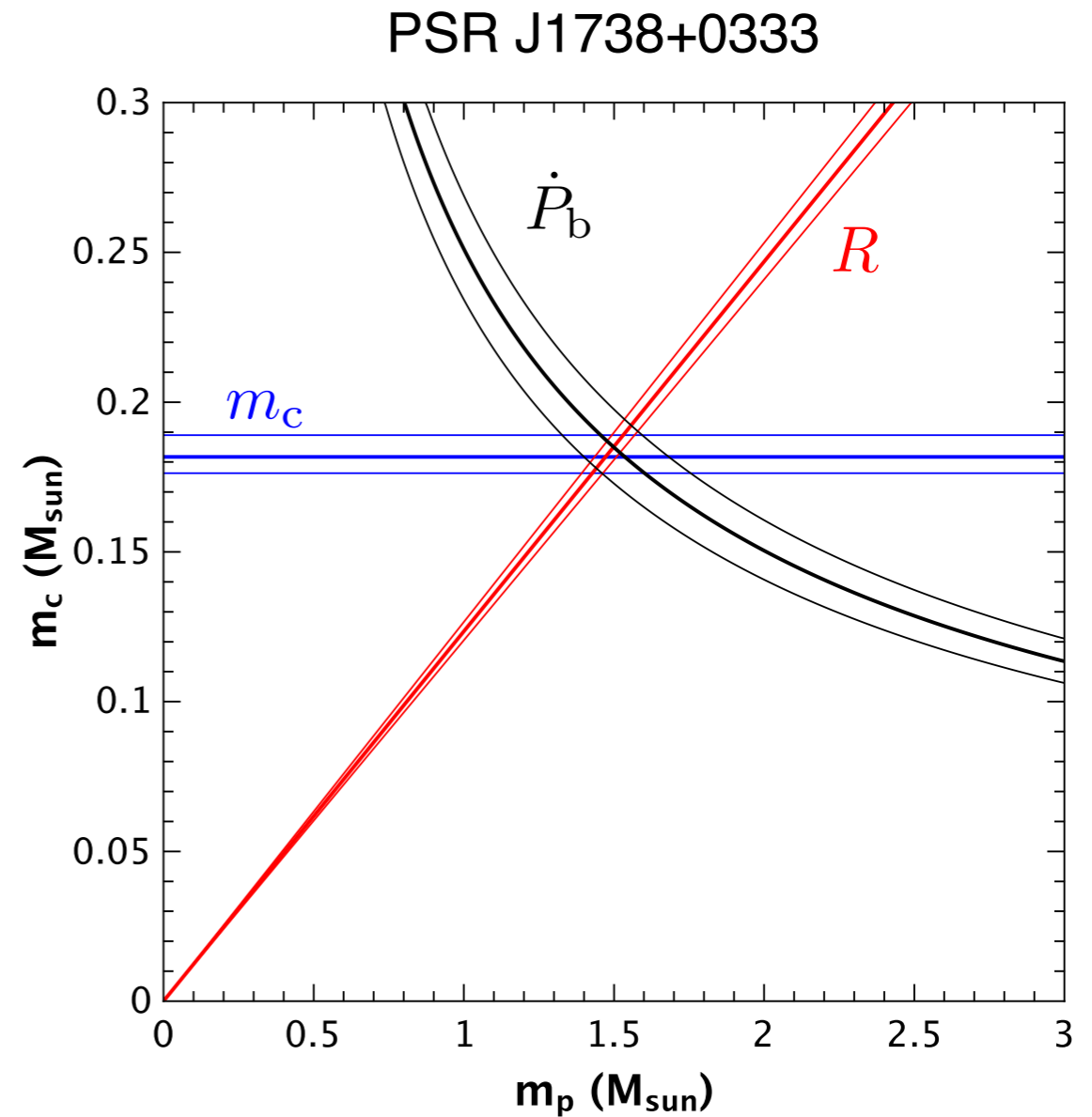


Double Pulsar agrees with GR significantly better than the Hulse-Taylor pulsar

PSRs J0348+0432 and J1738+0333



Antoniadis et al. 2013

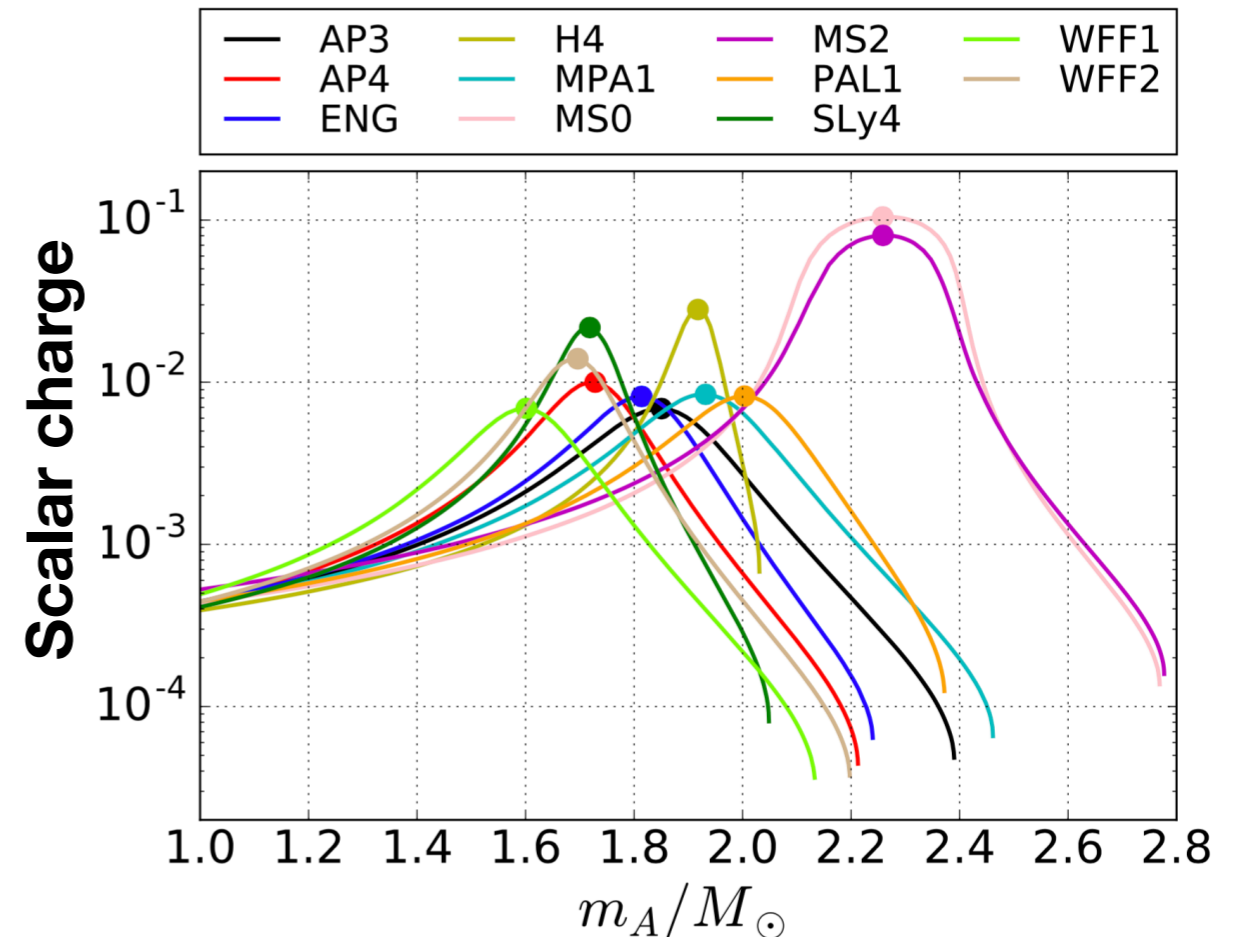
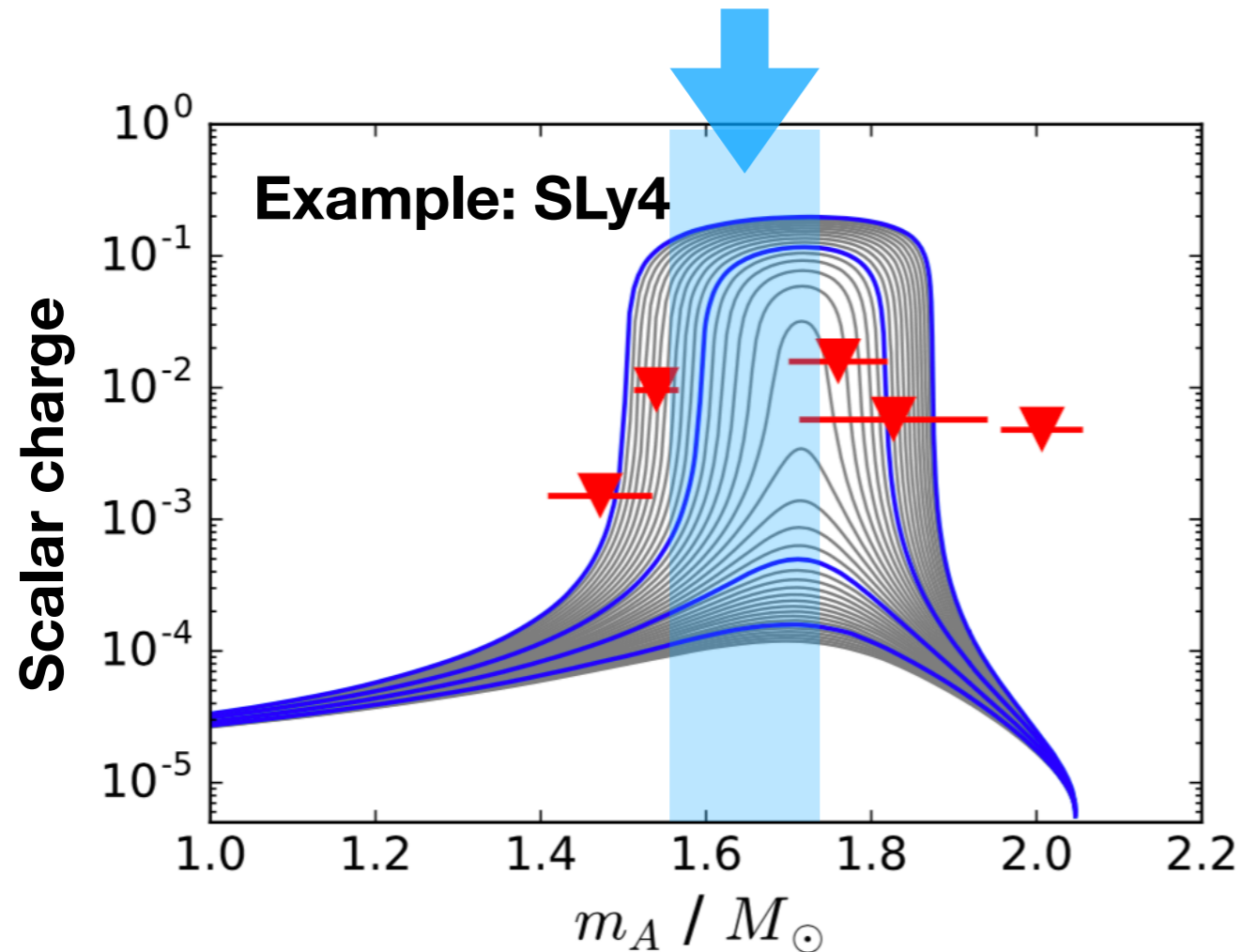


Freire et al. 2012

Due to their asymmetry, neutron-star white-dwarf systems provide stringent limits on *dipole radiation*

Combination of five NS-WDs

"Scalarization window"

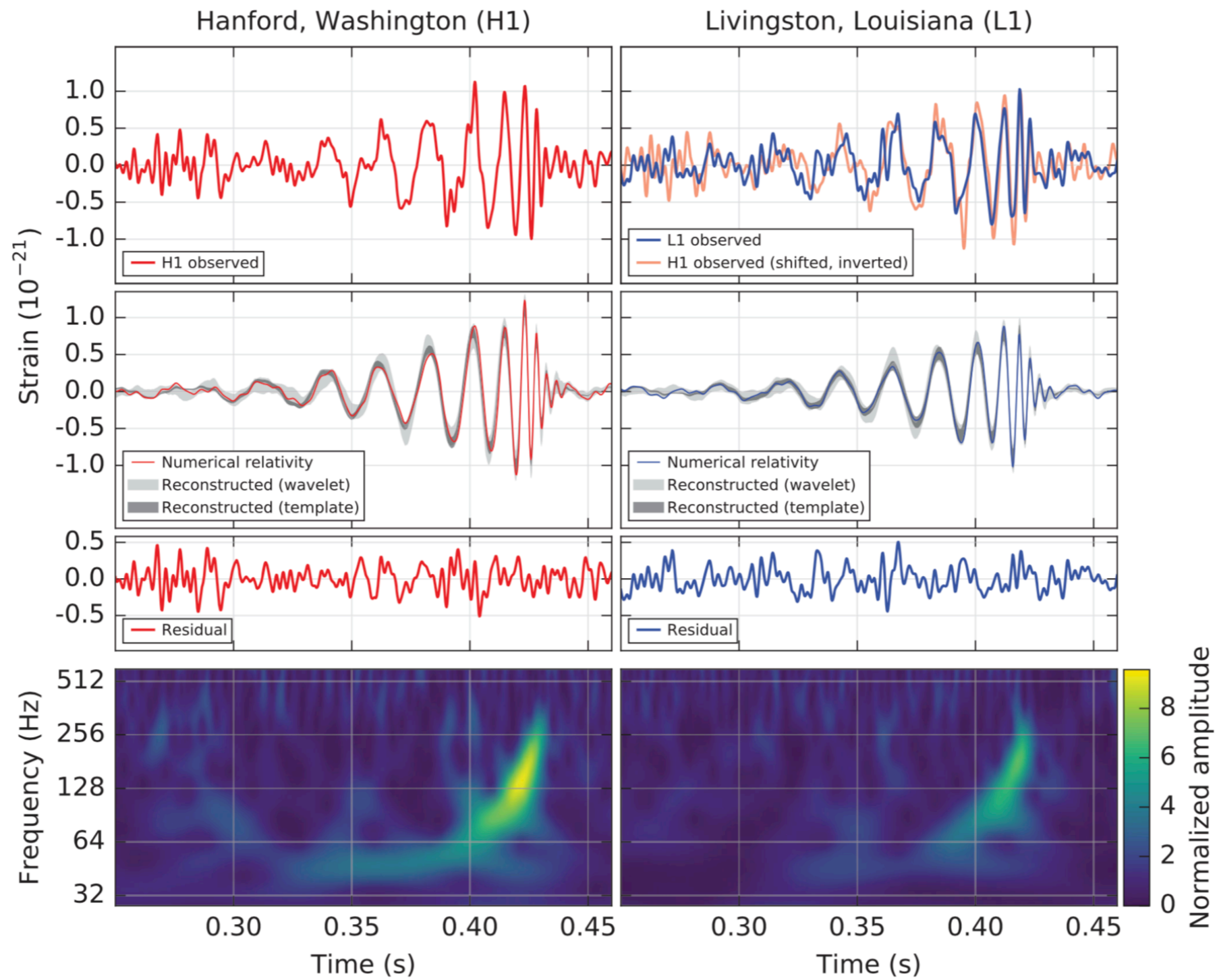


Strong-field effects could happen at different NS masses for different EOSs

Combining five best-timed NS-WD binaries put the best limits on a class of scalar-tensor theories for different EOSs

cf. Bastian's talk for EOSs

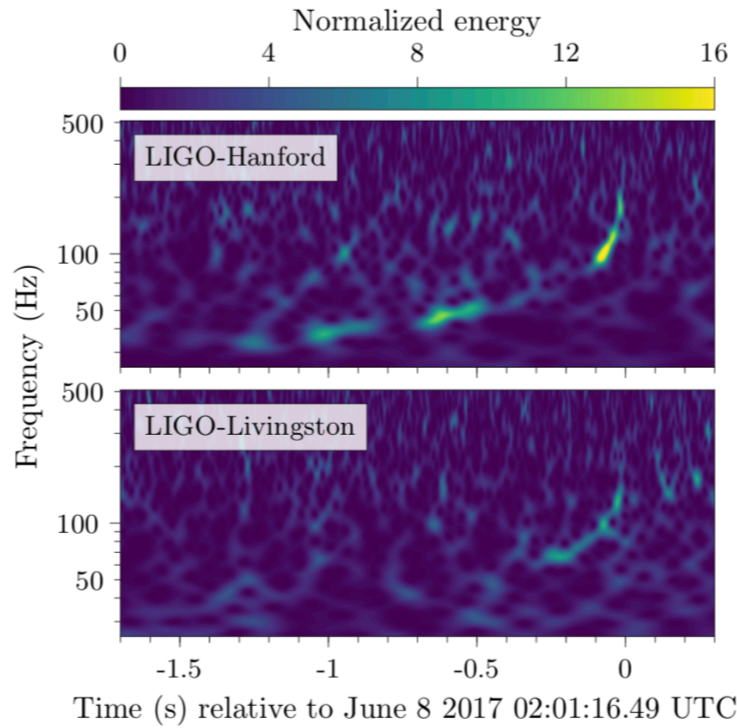
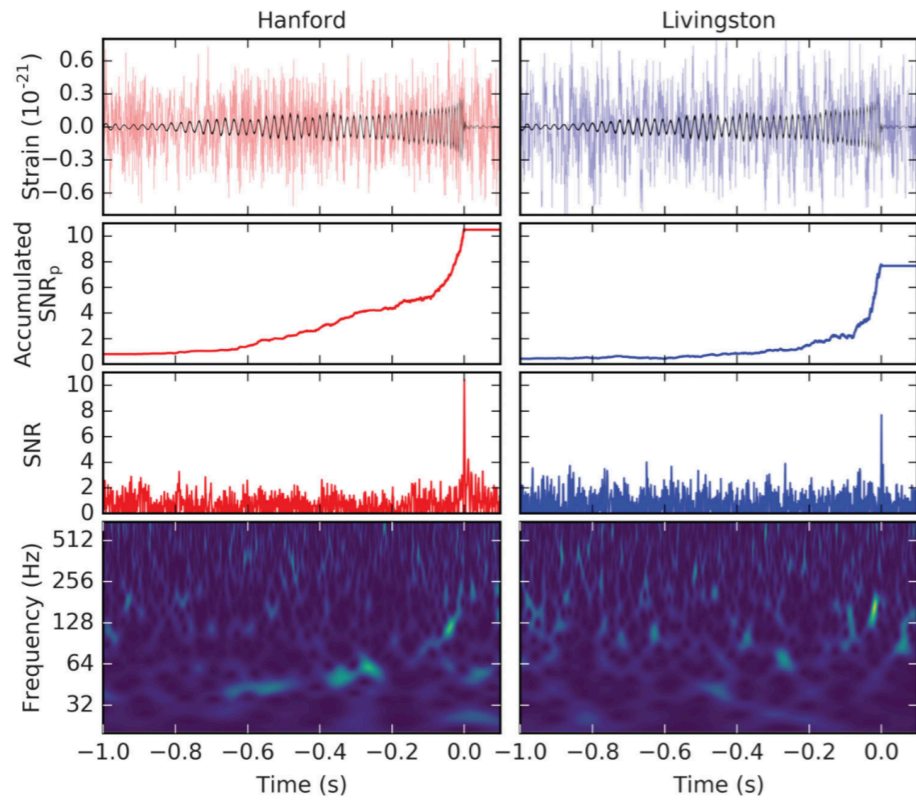
The first BBH merger: GW150914



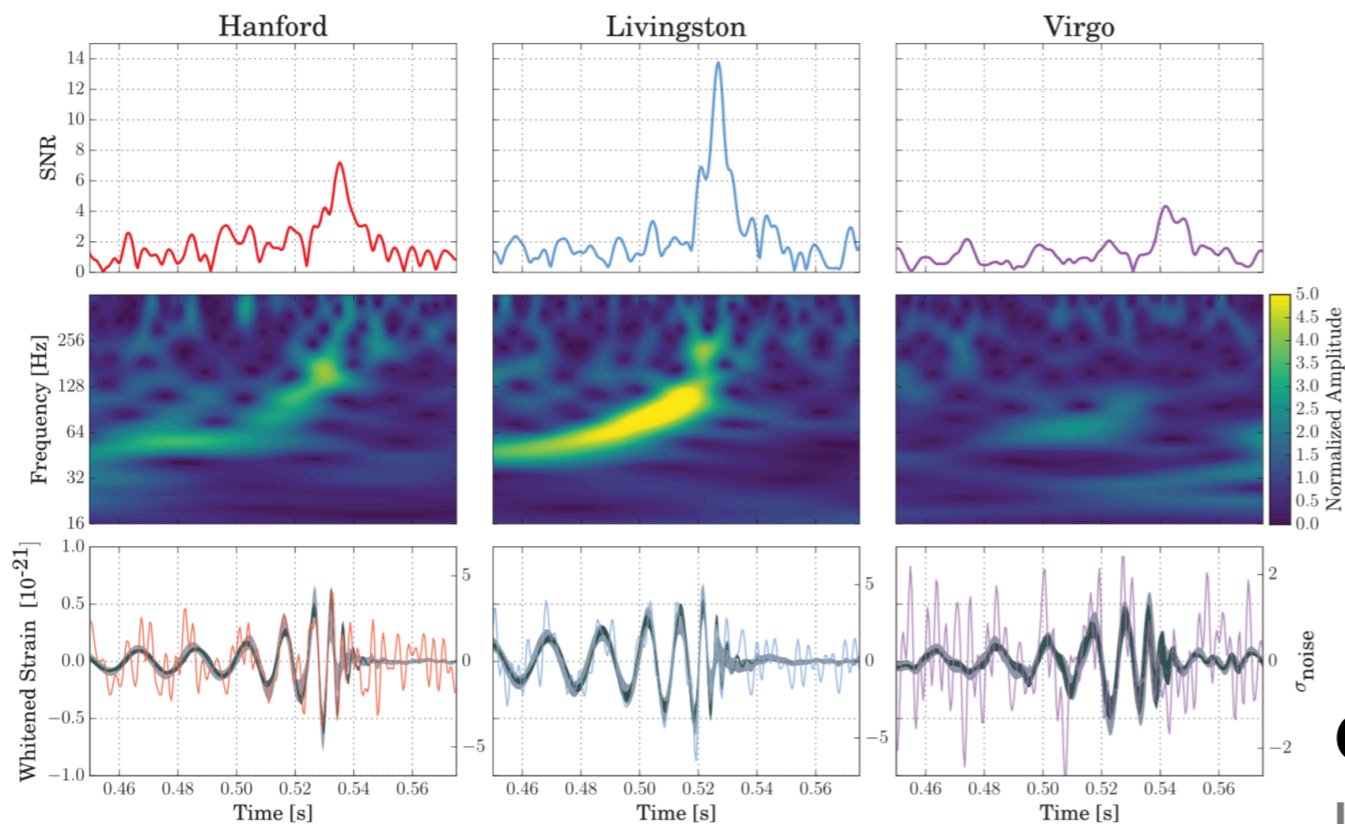
LVC 2016, PRL 116:061102

A zoo of BBH mergers

GW151226 LVC 2016, PRL 116:241103

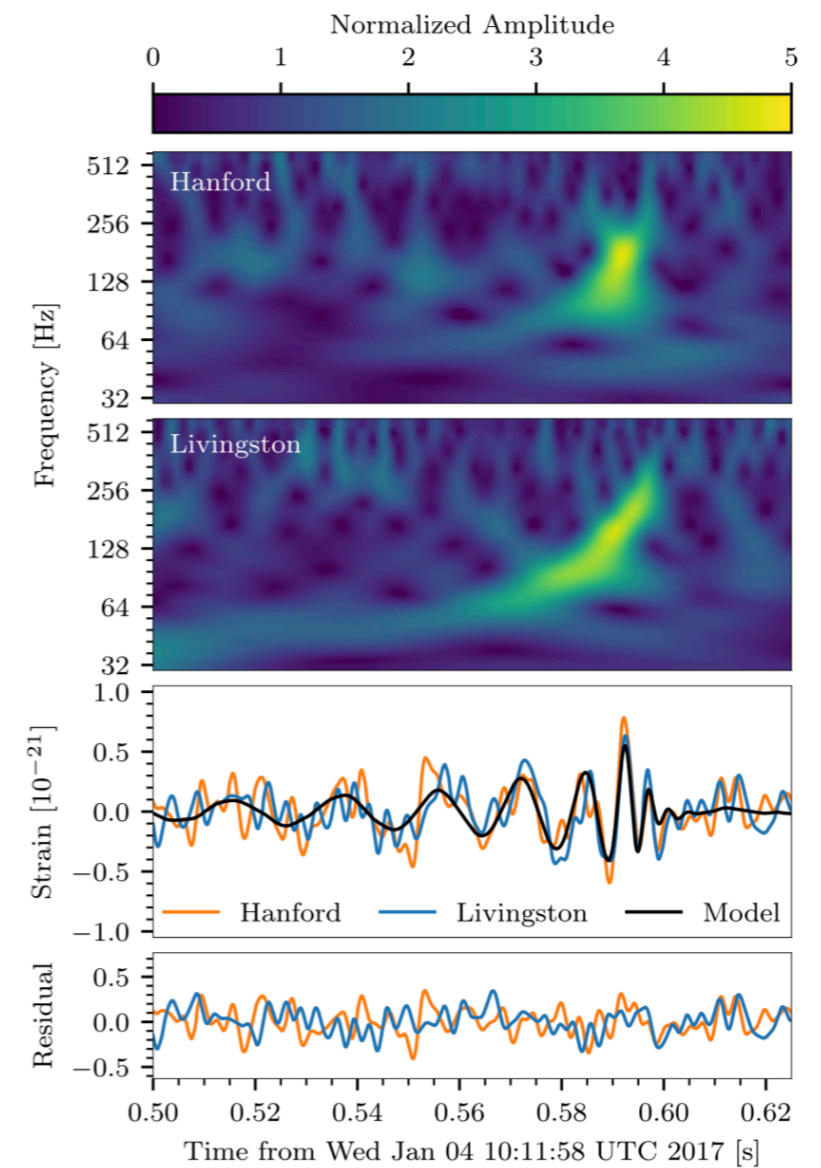


GW170608
LVC, arXiv:1711.05578



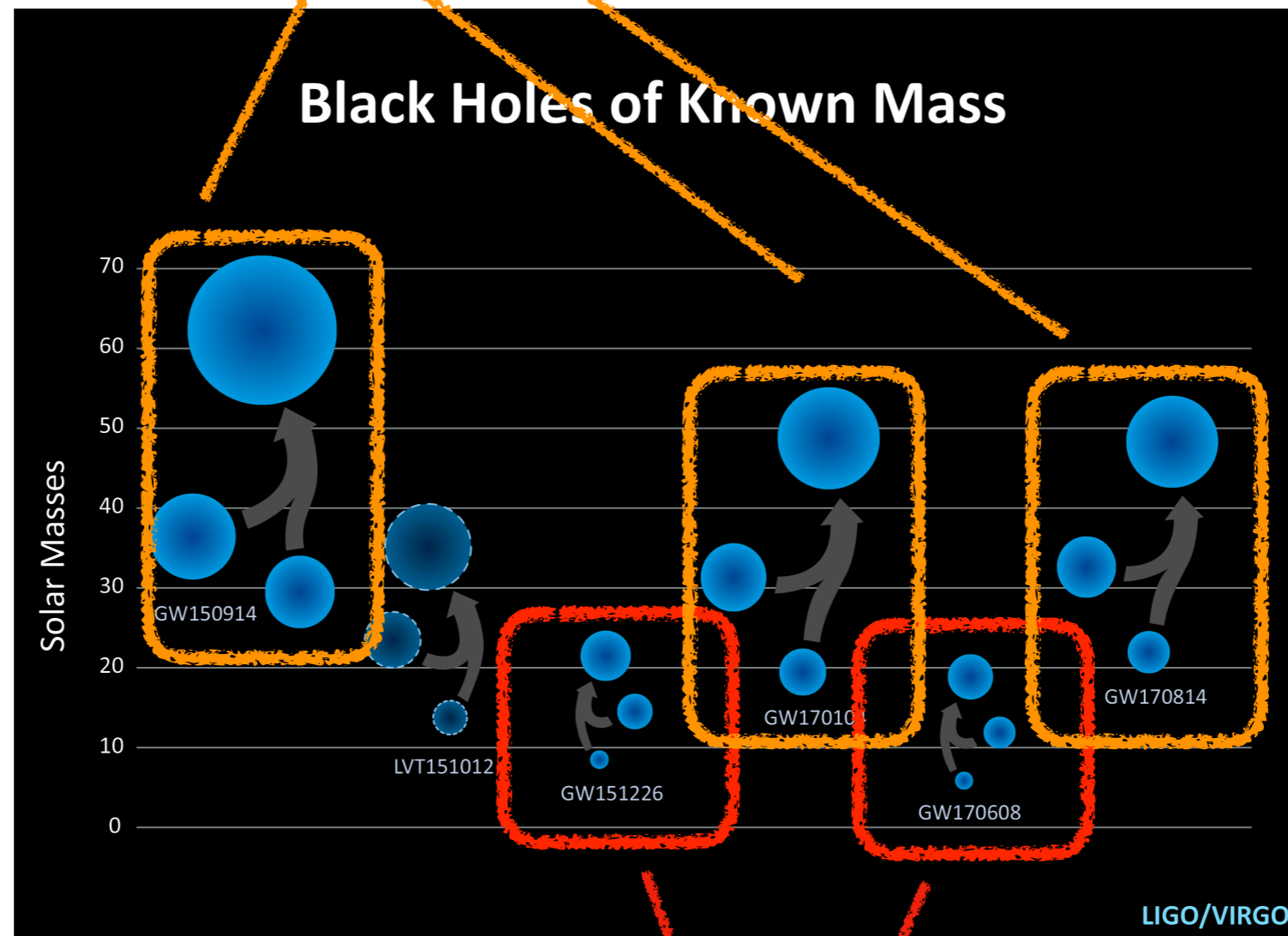
GW170814
LVC 2017, PRL 119:141101

GW170104
LVC 2017, PRL 118:221101



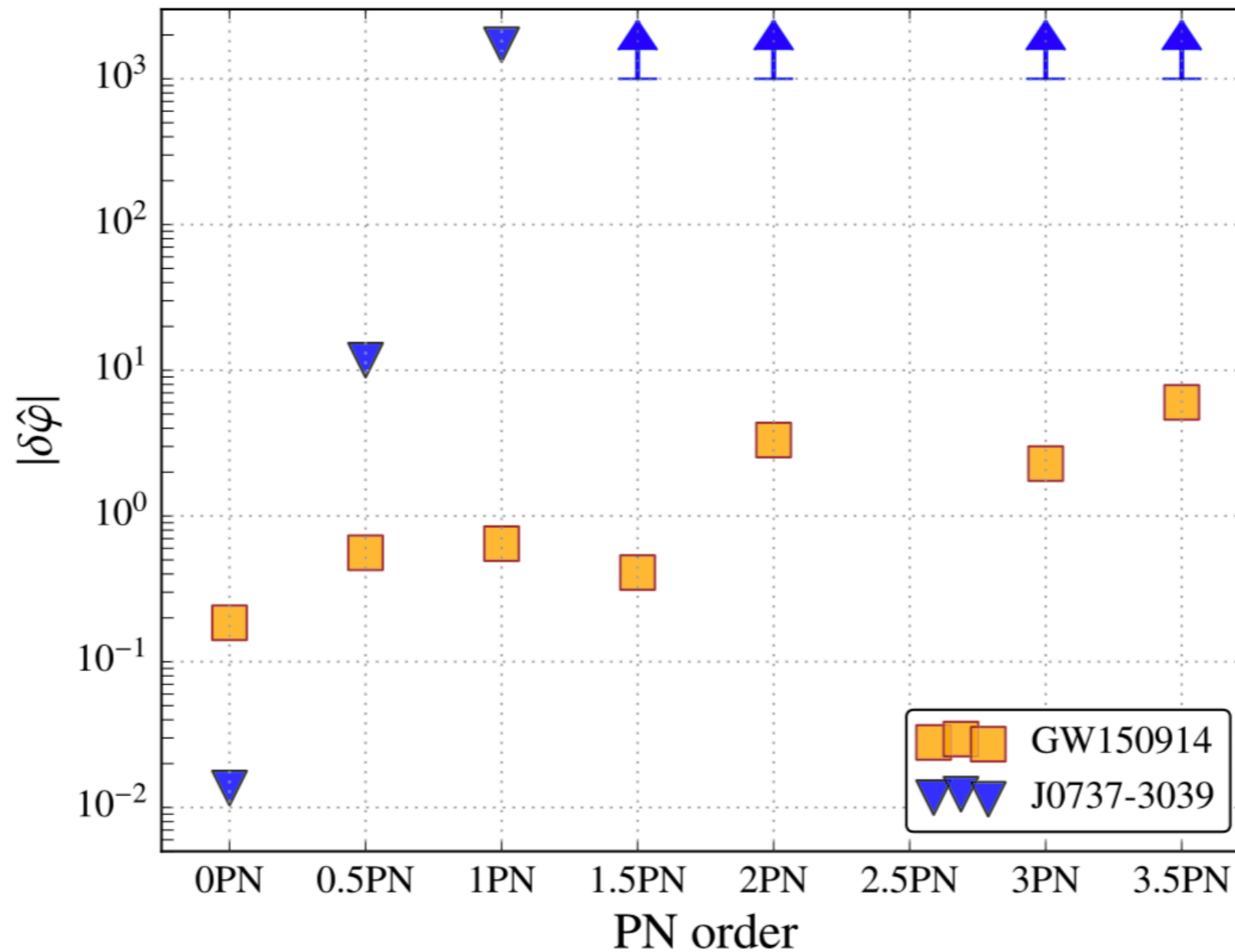
Masses of BBH mergers

Events with larger masses are in general better at probing merger and ringdown phases



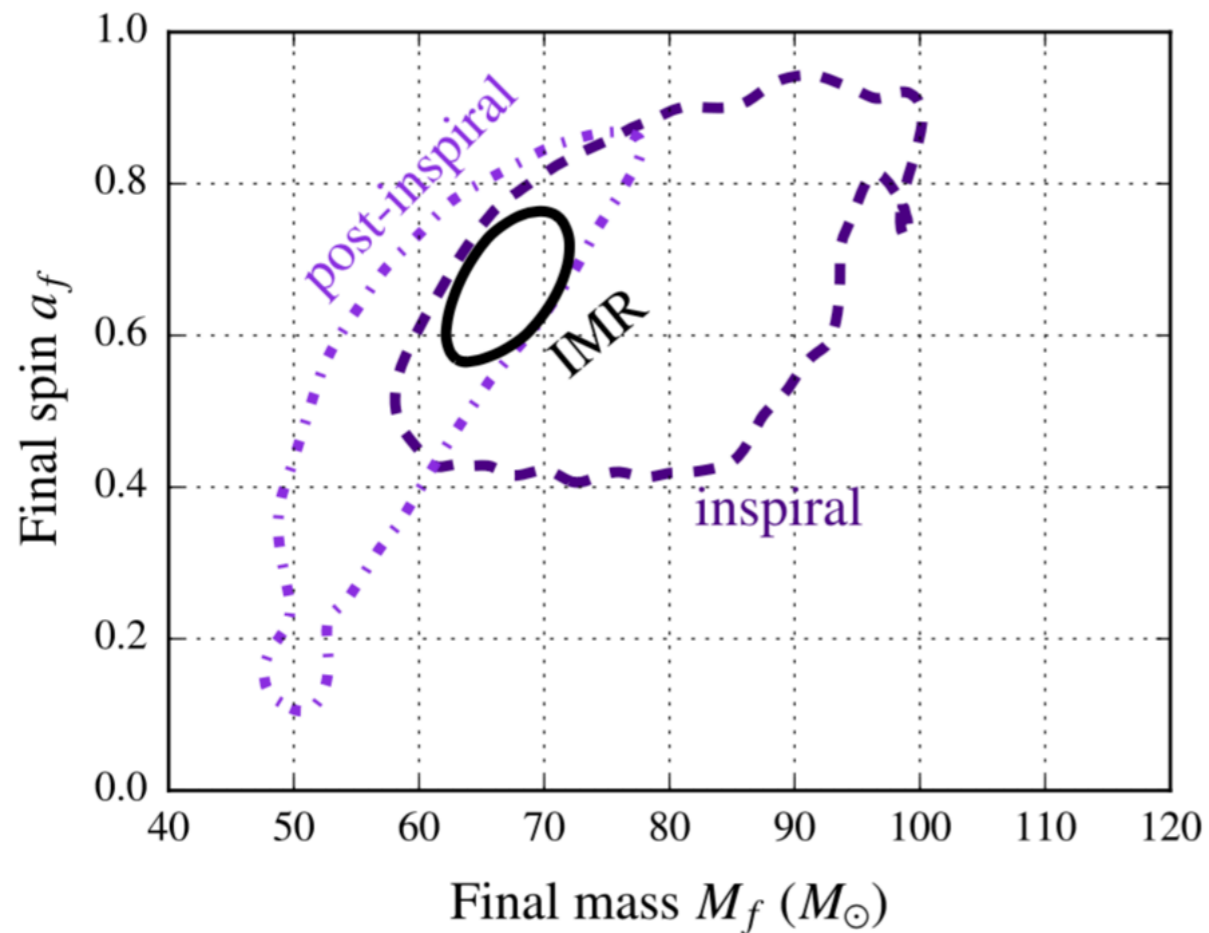
Events with smaller masses are in general better at probing lower-order post-Newtonian gravity

Tests of post-Newtonian dynamics



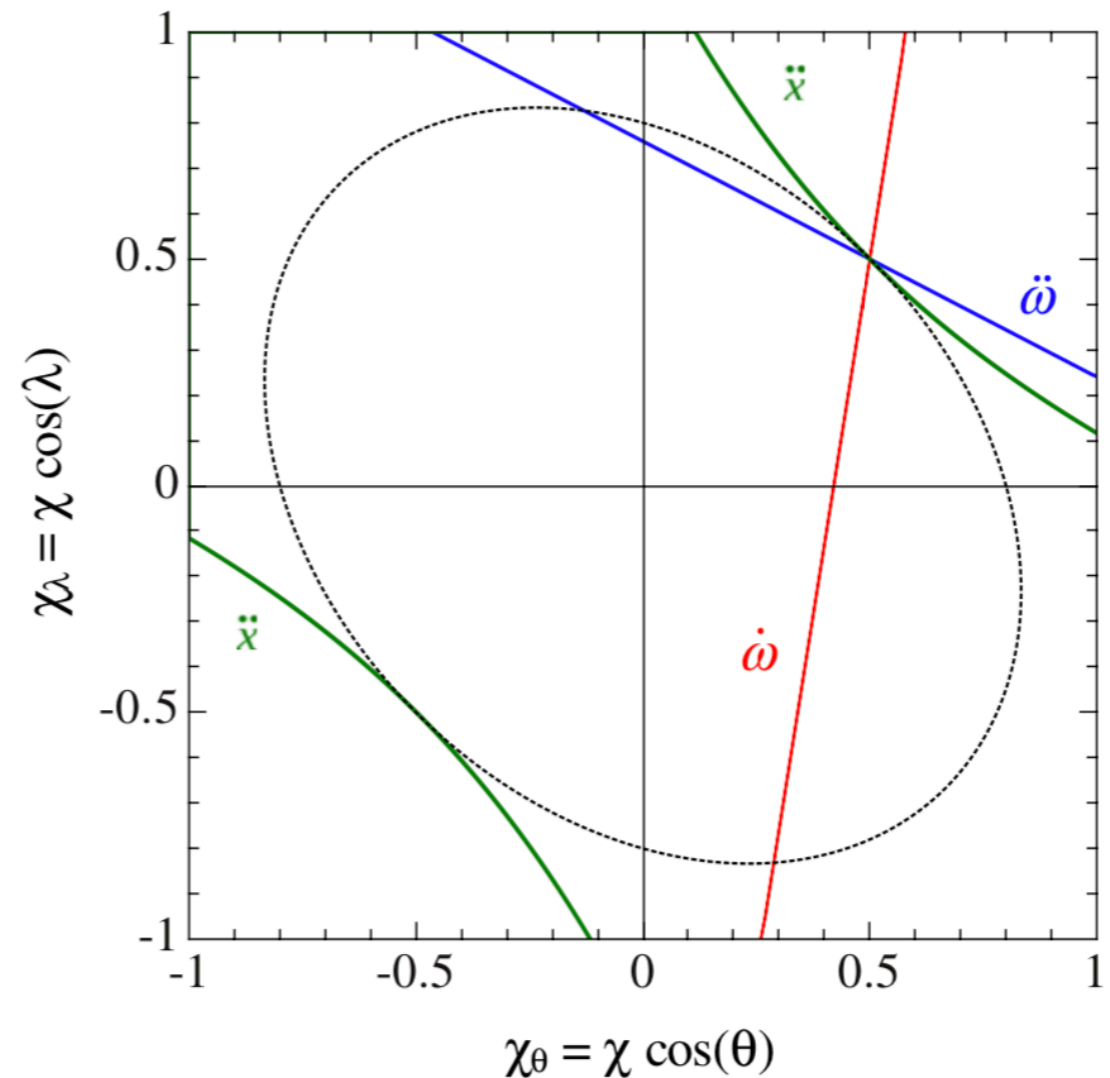
(using Double Pulsar results as of 2006)

IMR consistency and no-hair theorem



GW150914

Meidam et al. 2014, PRD 90:064009
LVC 2016, PRL 116:221101



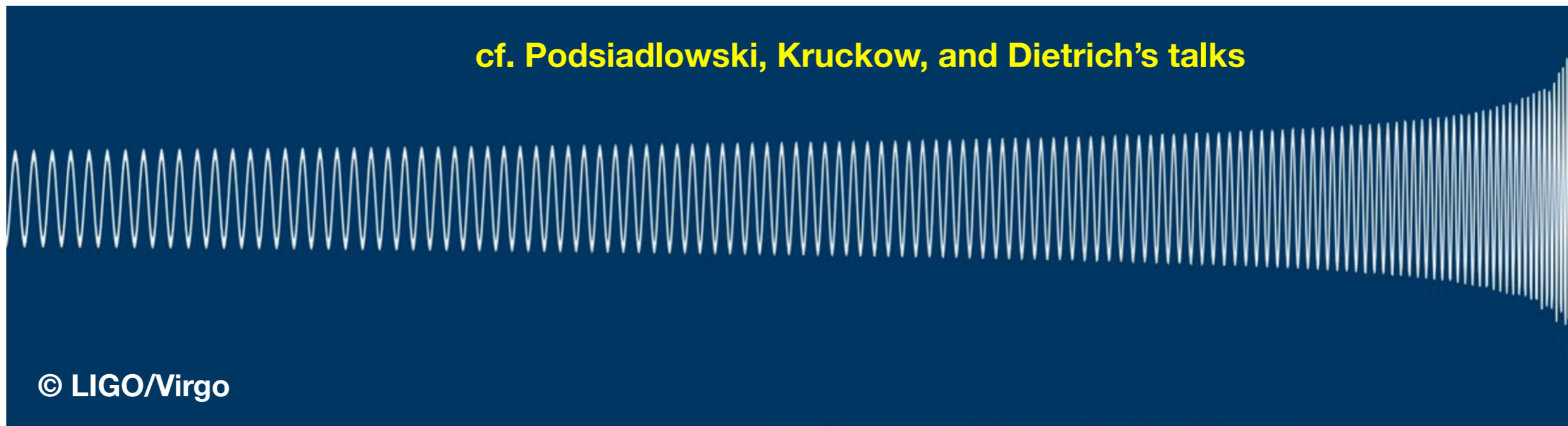
Pulsar around Sgr A*

Wex & Kopeikin 1999, ApJ 514:388
Liu et al. 2012, ApJ 747:1

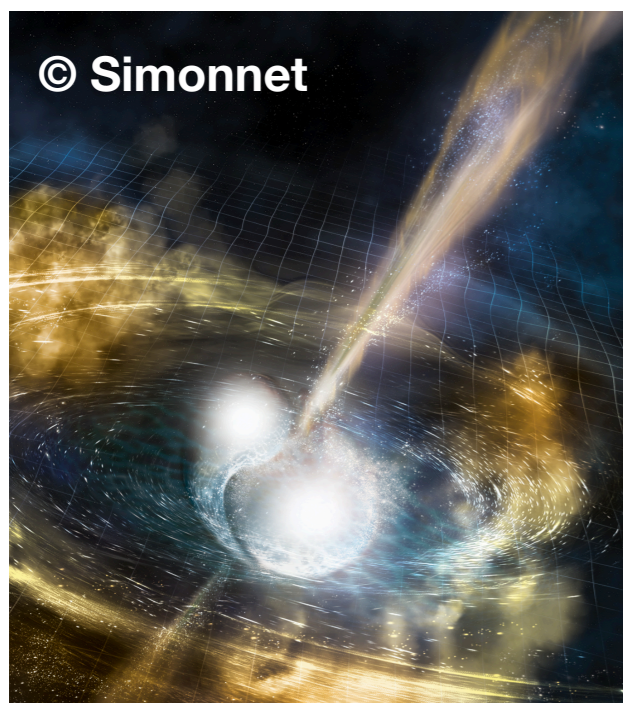
When higher-order quasinormal modes can be extracted with future GW detectors, no-hair theorem can be tested; it is also true when pulsars closely orbiting Sgr A* are discovered

GW170817: binary neutron stars

cf. Podsiadlowski, Kruckow, and Dietrich's talks

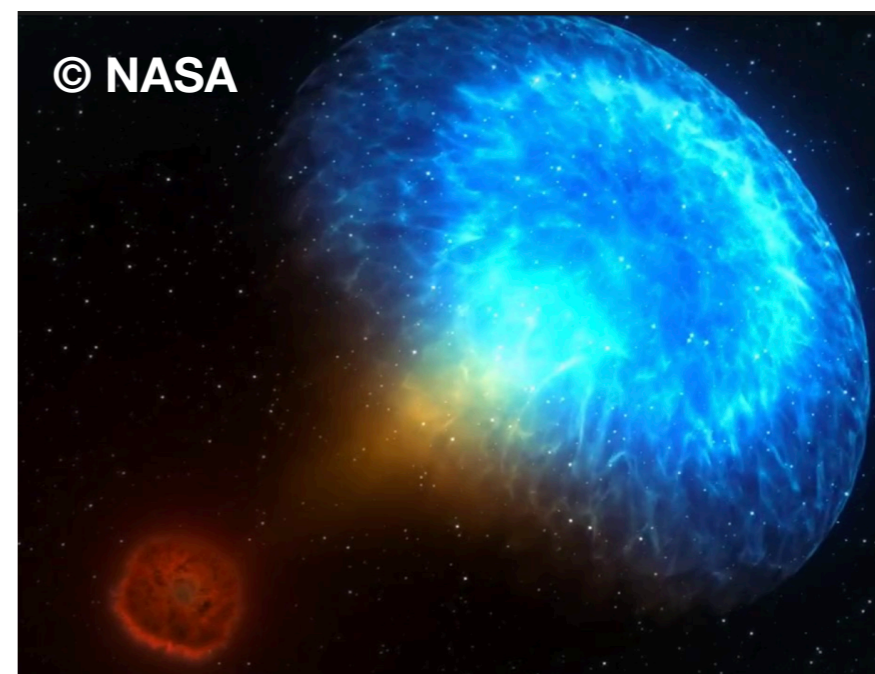


The long inspiral waveform could have imprints from the material of neutron stars \Rightarrow Testing Einstein's theory with matter

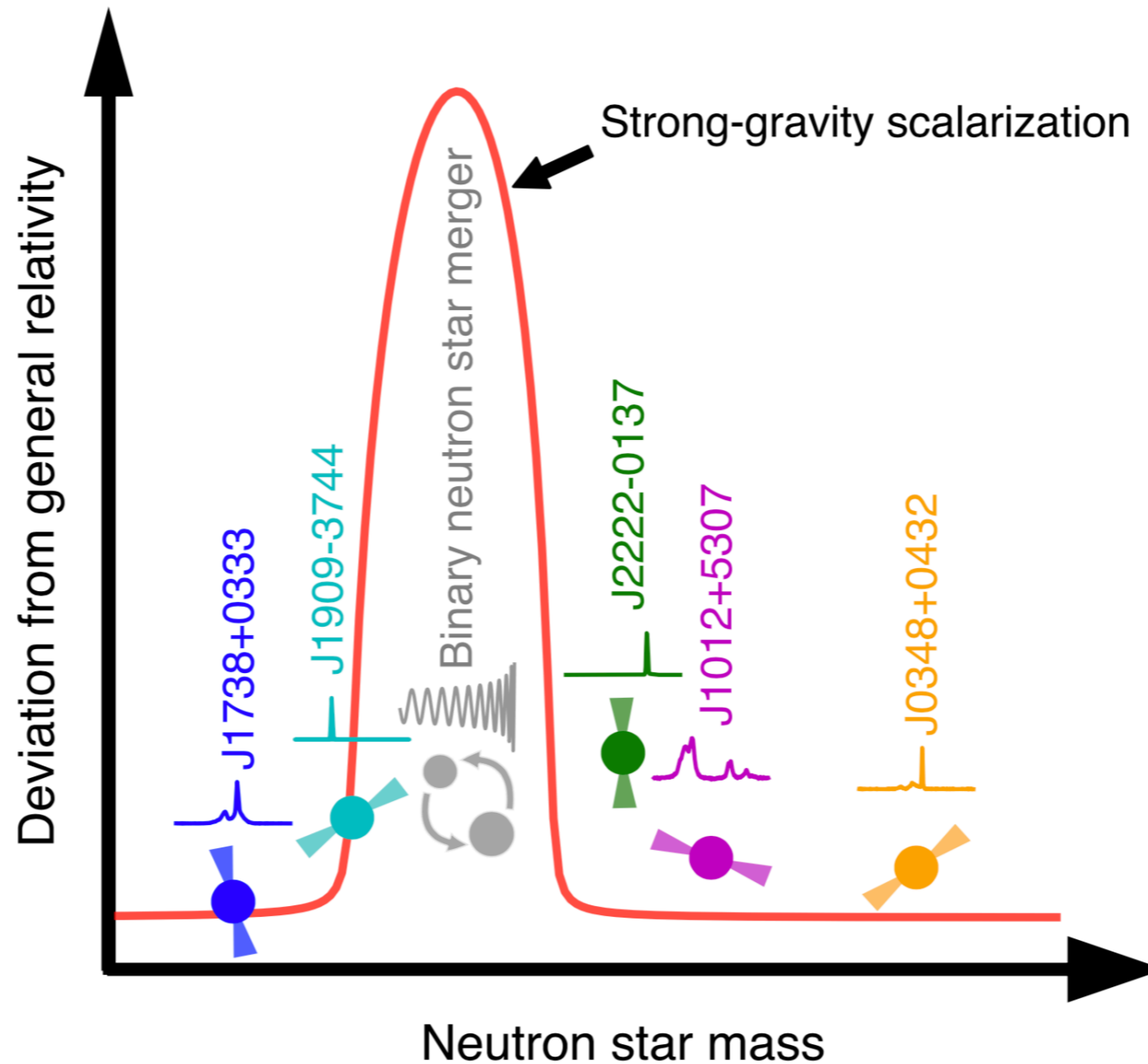


In addition, coincident observation with various electromagnetic signals probes the speed of gravity, implicating cosmology

LVC 2017, PRL 119:161101
LVC, Fermi, INTEGRAL 2017, ApJ 848:L13



Complementarity with pulsars and BNSs



If future GW detectors can observe BNSs with masses around 1.6–1.7 M_{sun} ,
a *scalarization window* will be closed

Summary

- ◆ **Despite various confirmation of GR in Solar System, in strong field gravitation might deviate from GR in a noticeable way**
- ◆ **Pulsars and GW detectors are superb laboratories to study SFG**
 - ✓ post-Newtonian dynamics
 - ✓ GW radiation (extra channels)
 - ✓ no-hair theorem
 - ✓ speed of gravity
 - ✓ ...
- ◆ **Pulsars and GW detectors are complementary in testing SFG in many ways, and a new horizon of experimental gravity is ahead of us**

Thank you!

